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Raman spectroscopic study of CO₂-NaCl-H₂O mixtures in synthetic fluid inclusions at high temperatures

JINYANG CHEN,^{1,*} HAIFEI ZHENG,² WANSHENG XIAO,¹ YISHAN ZENG,² and KENAN WENG¹ ¹Guangzhou Institute of Geochemistry, Chinese Academy of Sciences, Guangzhou 510640, China ²Department of Geology, Peking University, Beijing 100871, China

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Abstract—Mixtures of CO₂-NaCl-H₂O contained in synthetic fluid inclusions are studied by laser Raman spectroscopy at high temperatures. With increasing temperature, the band splitting (X) of v_1 - $2v_2$ diad of spectrum of CO₂ presents more variations, and the intensity ratios of the hot bands to the v_1 - $2v_2$ diad increase. For mixtures of gas phase rich in CO₂ and liquid phase rich in H₂O before homogenization, the strength of hydrogen bonding of water in the liquid phase decreases almost linearly with increasing temperature. When mixtures become homogeneous liquid phases, carbon dioxide content increases significantly as a result of the abrupt decrease in hydrogen bonds. Our results show that the hydrogen bonds change only slightly at higher temperatures above the homogeneous point, and a certain extent of the hydrogen bonds still remains at the highest temperature of 550°C of this work. The study is helpful to Raman spectroscopic analysis of natural fluid inclusions at high temperatures. *Copyright* © 2004 Elsevier Ltd

1. INTRODUCTION

Studies of hydrothermal fluids are important to understand phase equilibrium, dissolution, migration and deposition of mineral, abiogenic formation of oil and natural gas, geoelectric and geomagnetic discontinuity, and seismic wave (Duan et al., 2003; Xiao et al., 2001; Weng et al., 1999, 1997; Zheng et al., 1997; Seward et al., 1999). In addition, there are prospective applications to chemical engineering, material synthesis, disposal of hazardous waste, hydrometallurgy, desalinization of seawater, and geothermal exploitation (Savage et al., 1995; Yamanaka et al., 2001; Ziegler et al., 2001; Daimon et al., 2001; Mitton et al., 2001; Kritzer and Dinjus, 2001). In recent years, with the development of high-temperature and highpressure experimental technique and the applications of analytical instruments such as nuclear magnetic resonance (Hoffmann and Conradi, 1997; Matubayasi et al., 1997a, 1997b), X-ray diffraction (Ramos et al., 2000; Ohtaki et al., 1997), neutron diffraction (Okhulkov and Gorbaty, 2001; Ricci et al., 1998), microwave spectroscopy (Okada et al., 1997), X-ray absorption fine structure (Wallen et al., 1997, 1998; Seward et al., 1999; Bassett et al., 2000; Mayanovic et al., 2001; Anderson et al., 2002), Uv-vis (Suleimenov and Seward, 1997), infrared spectroscopy (Furutaka et al., 2001; Hu et al., 2000), and Raman spectroscopy (Frantz et al., 1993, 1994; Frantz, 1998; Fournier et al., 1998; Dubessy et al., 1999, 2001; Ebukuro et al., 1999; Walrafen et al., 1999; Ikushima et al., 1998a, 1998b; Carey and Korenowski, 1998; Carey et al., 1998), it is known that the strength of hydrogen bonding of water weakens with increasing temperature. At the critical point of pure water (374°C, 221 bar), the hydrogen bonds decrease abruptly. However, we have not sufficiently understood the nature of hydrogen bonding, and there are many arguments about whether they disappear or how many remain

CO₂-NaCl-H₂O is a typical geological fluid, and it is very important to geological processes in the deep earth. The properties of the hydrothermal fluid are closely related to the hydrogen bonding of water, for example, increasing temperature disrupts the hydrogen bonding and facilitates the insertion of CO₂, and therefore increases its solubility. Raman spectroscopy is probably the most important of the techniques that have been employed to obtain information relating to the intra- and intermolecular vibrational modes of water. The stretching vibration (v_2) of the Raman spectrum of water is sensitive to the hydrogen bonding (Frantz et al., 1993; Ikushima et al., 1998a); therefore, the aim of this paper is to characterize the modifications of the spectra of water of the CO2-NaCl-H2O system to give preliminary interpretations in terms of hydrogen bonds at high temperatures. In addition, the changes of spectra of CO_2 are examined in the prospective use of natural fluid inclusion analysis at high temperatures by micro-Raman spectroscopy.

Synthetic fluid inclusions give an ideal cell for high-temperature Raman spectroscopy. The scattering geometry used for synthetic fluid inclusions is similar to that of natural fluid inclusions; therefore, the results can be used for the analysis of natural fluid inclusions. Indeed, such an approach might be a good complement to classical optical cells for its unique geometry.

2. MATERIAL AND METHODS

Fluid inclusions were synthesized by healing fractures in natural quartz cores at various temperatures and pressures using the classical technique (Sterner and Bodnar, 1984; Bodnar and Sterner, 1985; Frost

at higher temperatures above the critical point (Postorino et al., 1993; Tromp et al., 1994; Hoffmann and Conradi, 1997; Gorbaty and Kalinichev, 1995; Matubayasi et al., 1997a, 1997b; Soper et al., 1997; Bellissent-Funel, 2001; Bellissent-Funel and Tassaing, 1997; Marti, 1999, 2000; Marti et al., 1996; Yamaguchi et al., 2001; Mayanovic et al., 2001). Furthermore, the effect of other species such as nondipolar molecules (CH₄, N₂, CO₂) on the hydrogen bonding should be studied in detail.

^{*} Author to whom correspondence should be addressed (chenjy@ gig.ac.cn).

Sample T (°C) P (MPa) 1 2 3 4 $10\% \text{ CO}_2 + \text{H}_2\text{O}$ 595.0-597.2 COHG 650 100 595.8 (8) Gas COHL 600 120 366.0-368.9 367.8 (7) $5\% \text{ CO}_2 + \text{H}_2\text{O}$ Liquid $5\% \text{ NaCl} + 5\% \text{ CO}_2 + \text{H}_2\text{O}$ SCOHL 600 120 399.1-402.0 400.5 (6) Liquid SCOHG 5% NaCl + 10% CO₂ + H₂O 600 100 529.1-531.7 530.2 (7) Gas

Table 1. Forming temperatures and pressures, homogenization temperatures, and compositions of synthetic fluid inclusions.

(1) Range of measured homogenization temperatures (°C).

(2) Average homogenization temperature (Th) (°C) and number of measurements (in parenthesis).

(3) Composition in weight percent.

(4) Phase of homogenization.

and Wood, 1997; Shmulovich and Graham, 1999; Bakker and Diamond, 2000; Schmidt and Bodnar, 2000). The CO_2 was obtained from the decomposition of $Ag_2C_2O_4$. After the synthetic experiments, the quartz cores were cut into approximately 1-mm-thick disks, polished on both sides, and examined with a petrographic microscope. Almost all fluid inclusions distribute along the fractures of the quartz cores, and the liquid–gas ratios of every sample are almost equal, which indicates that fluid inclusions of every sample trapped a homogeneous fluid with the same composition and density as the parent solution (Sterner and Bodnar, 1984). The homogeneous temperatures (Th) of these samples were measured in Leitz 1350 stage. The forming temperatures and pressures, homogenization temperatures, and compositions of synthetic fluid inclusions are given in Table 1. Because the ranges of measured Th are less than 3°C, the results also confirm that the compositions and densities of the fluid inclusions of every sample are uniform.

Raman spectra were obtained by a Raman micro-spectrometer (RENISHAW System RM-1000, Renishaw group, Gloucestershire, United Kingdom) equipped with a heating stage (Linkam TM93). The measurements use an ionized Ar⁺ laser, $\lambda_0 = 514.5$ nm, 50 mW, entrance slit 50 nm, collection time 30 s, and range 1000~4000 cm⁻¹. Cold water flowing through a copper tube, wound around the long-working-distance objective (×50, Olympus), allowed acquisition of spectra to 600°C without damaging the objective. Figure 1 clearly shows that mixtures of gas phase and liquid phase in fluid inclusions become homogeneous gradually during heating.

3. RESULTS

A typical Raman spectrum obtained from a synthetic fluid inclusion is illustrated in Figure 2a (SCOHG, gas phase: 330° C). At low wavenumbers, the bands of the quartz host crystal are identified, and they do not overlap with the bands of carbon dioxide and water.

Carbon dioxide exhibits two bands around 1380 and 1278 cm⁻¹ as results of Fermi resonance of v_1 and $2v_2$ (Wienecke et al., 1986; Rosso and Bodnar, 1995), and two obvious shoulder peaks called hot band near Fermi diad (Figs. 2 and 3). Before homogenization, the gas phase rich in CO₂ yields two strong bands of carbon dioxide, whereas the liquid phase rich in H₂O gives two weak bands of carbon dioxide (Fig. 2b). In this study we discuss the strong bands of carbon dioxide obtained in gas phases only.

Figure 3 represents spectra of CO_2 in the gas phase of COHG (10%CO₂ + H₂O, Th (G) = 596°C), and it shows that the

80°C 580°C 380°C 480°C COHG 130°C 230°C 330°C 368°C COHL 330°C 395°C 400°C 80°C SCOHL 80°C 230°C 380°C 530°C SCOHG 20 µ m

Fig. 1. Variation of mixtures in fluid inclusions during heating.

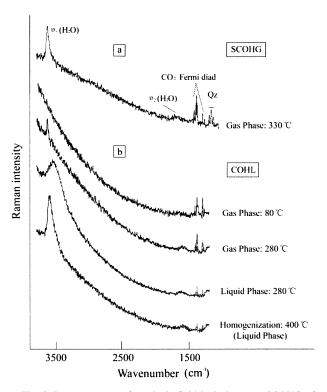


Fig. 2. Raman spectra of synthetic fluid inclusions. (a) SCOHG of the gas phase at 330° C; (b) COHL of different phases and temperatures.

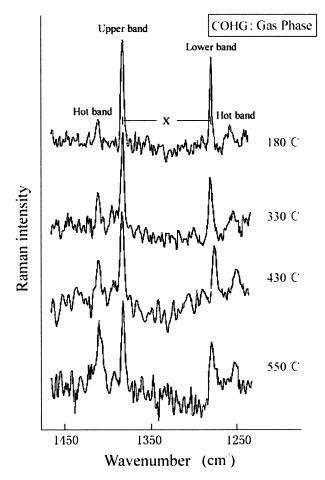


Fig. 3. Raman spectra of CO_2 in the gas phase of COHG at different temperatures. X is the Fermi diad splitting between the upper and lower bands.

intensity ratios of the hot bands to the Fermi diad increase with increasing temperature. For the Fermi diad splitting (X) between the upper and lower bands of CO_2 in the gas phases of COHG and SCOHG (5% NaCl + 10% CO_2 + H₂O, Th (G) = 530°C), Figure 4 shows that the variations of x values are small at lower temperatures, whereas, with increasing temperature, they present more and more variations.

Water exhibits two Raman bands assigned to internal motions, and they are the bending mode of water (v_2) around 1600 cm^{-1} and a large massif assigned to the stretching vibration (v_s) in 3000~3700 cm⁻¹ (Walrafen, 1964, 1967; Ratcllffe and Irish, 1982; Frantz et al., 1993) (Fig. 2). Figure 2b shows that before homogenization, the gas phase gives very weak Raman bands of water, whereas the bands of the liquid phase are very strong, especially for the stretching band. The bending mode of water (v_2) is weak and has not been studied in this work. As for the very strong spectra of the stretching vibrational mode of water in the liquid phase, with increasing temperature, the massif becomes sharper and sharper, the frequency at maximum intensity (v_s) shifts to higher wavenumbers, and the full width at half maximum intensity (FWHM) decreases (Fig. 5). When compared with spectra of carbon dioxide, those for the stretching band of water are the ones that exhibit the highest

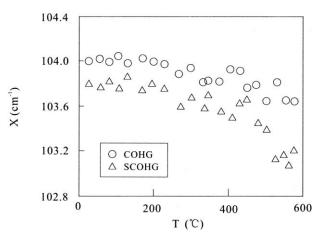


Fig. 4. The Fermi diad splitting (X) between the upper and lower bands of CO_2 in the gas phases of COHG and SCOHG with temperature.

variations in frequency at maximum intensity (v_s), FWHM, and shape.

Figure 6 illustrates the changes of the v_s and FWHM values of the stretching band of water in the liquid phases for COHL (5% CO₂ + H₂O, Th (L) = 368°C) and SCOHL (5% NaCl + 5% CO₂ + H₂O, Th (L) = 400°C) with temperature. It shows that the v_s and FWHM values increase and decrease almost

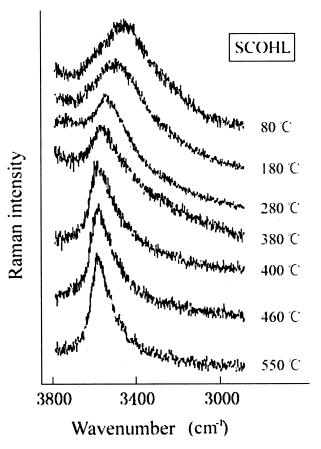


Fig. 5. Raman spectra of the stretching band of water in the liquid phase of SCOHL at different temperatures.

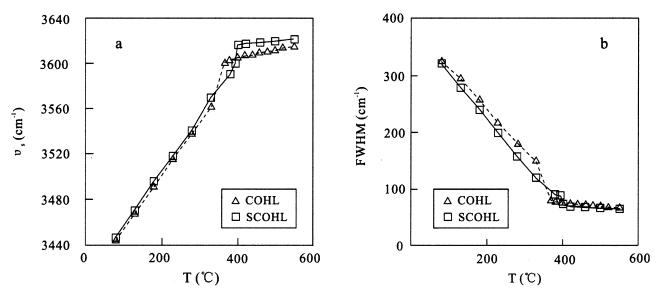


Fig. 6. Variation of v_s and FWHM of the stretching band of water in the liquid phases of COHL and SCOHL with temperature. (a) v_s ; (b) FWHM.

linearly with increasing temperature before homogenization, respectively; near the Th, they both change abruptly; whereas, for the homogeneous liquid phases, the variations of the v_s with temperature are small.

4. DISCUSSION

4.1. Raman Vibration of CO₂

The Raman spectrum for a homogeneous, free CO_2 phase shows that the Fermi resonance splitting (X) between the upper and lower bands is almost linearly related to the pressure at ambient temperature (Rosso and Bodnar, 1995); whereas most CO_2 in natural fluid inclusions are not homogeneous, free phases, and often mix with aqueous solutions. Therefore, the spectra of carbon dioxide in the gas phase are characterized in this study for COHG (10% $CO_2 + H_2O$, Th (G) = 596°C) and SCOHG (5% NaCl + 10% $CO_2 + H_2O$, Th (G) = 530°C).

At lower temperatures, water content in gas phases is comparatively low, and its effect can be ignored; therefore, the temperature is regarded to cause the main effect on Raman spectra of carbon dioxide. From the fact that the X values in the gas phases of COHG and SCOHG show small changes with increasing temperature at lower temperatures (Fig. 4), it can be inferred that the temperature has little effect on X at this condition. Therefore, for multiphase, pure CO_2 in fluid inclusions, the densities might be determined through heating to homogeneous gases by Raman spectroscopy as Rosso and Bodnar (1995) did at ambient temperature. At higher temperatures, the X values present very obvious variations; the increase in water content in the gas phases and a nonlinear temperature effect might account for this phenomenon.

The Th of SCOHG of 530°C is lower than that of COHG of 596°C, indicating that at the same temperature the water content in the gas phase of SCOHG is more than that of COHG; as a result, the curve of X values as a function of temperature for SCOHG is lower and with more variations than that for

SCOHG (Fig. 4). When SCOHG becomes a homogeneous gas phase, the abrupt change of X values might mainly result from the significant increase of water and NaCl contents in the gas phase.

Hot bands near the Fermi diad of carbon dioxide are the result of two effects (Rosso and Bodnar, 1995; Dubessy et al., 1999): 1) the anharmonicity of the vibrational levels producing different energetic differences between two successive vibrational levels (E[v = 2] – E[v = 1] \neq E[v = 1] – E[v = 0]); and 2) the population of the excited vibrational levels increases with temperature according to the Boltzman law. Owing to the fact that Fermi resonance doubles the number of hot bands, there are two hot bands near the fundamental bands, at 1380 and 1278 cm^{-1} . With increasing temperature, the population of thermally excited vibrational levels increases, therefore, the intensity ratio of the hot bands to the Fermi diad increases (Fig. 3), which indicates that the intensity ratio might be used to estimate temperature, for example, the temperature of an externally heated diamond anvil cell (DAC) can be calibrated by this method.

4.2. Hydrogen Bonding of Water

Owing to the strong H₂O-H₂O interactions and the partial weakening of the covalent O-H stretching due to the existence of an O-H-O bridge, the Raman spectrum of the stretching band of water exhibits the highest variations in both frequency at maximum intensity (v_s) and FWHM. Therefore, we can obtain the changes of the strength of the hydrogen bonding by the variations of the v_s and FWHM (Frantz et al., 1993; Ikushima et al., 1998a). As a result of increasing temperature, the bands of stretching vibration of water become sharper and sharper and the v_s shift to higher wavenumbers (Fig. 5), indicating that the hydrogen bonds decrease with increasing temperature. For the spectra of water in the liquid phases before homogenization, Figure 6 shows that the v_s and FWHM values increase and

decrease almost linearly with increasing temperature, respectively, suggesting that almost linear decreases of hydrogen bonds result from increasing temperature. There are liquid phase and gas phase coexisting before homogenization, so the changes of the pressures in the inclusions are very slight with temperature, and it is assumed to be isobaric, indicating that increasing temperature results in almost linear decreases of the hydrogen bonds on this condition.

Comparing the spectra of the stretching bands of water in the liquid phase of SCOHL (5% NaCl + 5% CO₂ + H₂O, Th (L) = 400°C) with that of COHL (5% CO₂ + H₂O, Th (L) = 368°C before homogenization, the v_s and FWHM values of SCOHL lie above and below the corresponding value of COHL, respectively (Fig. 6). The liquid phase of SCOHL contains 5% NaCl, whereas that of COHL contains no NaCl; thus, NaCl might account for the difference of the spectra. During heating of the fluid mixtures in COHL and SCOHL, and as the systems approach the homogenization temperatures, the v_s and FWHM values change abruptly (Fig. 6), indicating a rapid decrease in the hydrogen bonds caused by the significant increase of CO₂ content in the liquid phase during homogenization.

Differing from the great changes of the stretching bands of water in the liquid phases below Th, with increasing temperature, the shape of these bands for the homogeneous liquid phases presents less change, and the variations of the v_s and FWHM values are small (Figs. 5 and 6). It can be assumed that the fluids are isochoric above Th, the changes of the spectra on this condition are slight with increasing temperature, indicating that the interactions of water molecules are slightly affected and the strength of the hydrogen bonding remains almost constant here.

At the maximum temperature of 550°C of this study, the v_s values of COHL and SCOHL are 3614 and 3621 cm⁻¹ respectively. Compared with the v_s value of 3756 cm⁻¹ for isolated monomer water (Frantz et al., 1993), these lower values indicates that the interactions of water molecules still exist and a certain extent of the hydrogen bonding network remains at 550°C for these fluids.

5. CONCLUSION

At lower temperatures, the variation of Fermi resonance splittings (X) of Raman spectra for carbon dioxide is small only with increasing temperature, so the densities of multiphase, pure CO_2 in natural inclusions might be determined using Raman spectroscopy by heating them to a homogeneous phase as Rosso and Bodnar (1995) did at ambient temperature. The intensity ratios of the hot bands to the Fermi diad increase with temperature, and this provides a useful method for the determination of temperature by Raman spectroscopy.

Increasing temperature results in almost a linear decrease of hydrogen bonds of water at isobaric condition. CO_2 and NaCl present have some effect in destroying hydrogen bonding, and the molecular interaction mechanism of CO_2 -NaCl-H₂O mixtures should be studied further to reveal the details of the changes of the hydrogen bonding. Our results show that the hydrogen bonds remain close to the temperature of 550°C used in the present study, and therefore, the hydrogen bonds of water

at the temperatures above the critical temperature of water (374°C) can not be ignored.

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