

# Shear wave velocity, seismic attenuation, and thermal structure of the continental upper mantle

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## SUMMARY

Seismic velocity and attenuation anomalies in the mantle are commonly interpreted in terms of temperature variations on the basis of laboratory studies of elastic and anelastic properties of rocks. In order to evaluate the relative contributions of thermal and non-thermal effects on anomalies of attenuation of seismic shear waves,  $Q_s^{-1}$ , and seismic velocity,  $V_s$ , we compare global maps of the thermal structure of the continental upper mantle with global  $Q_s^{-1}$  and  $V_s$  maps as determined from Rayleigh waves at periods between 40 and 150 s. We limit the comparison to three continental mantle depths (50, 100 and 150 km), where model resolution is relatively high.

The available data set does not indicate that, at a global scale, seismic anomalies in the upper mantle are controlled solely by temperature variations. Continental maps have correlation coefficients of  $<0.56$  between  $V_s$  and  $T$  and of  $<0.47$  between  $Q_s$  and  $T$  at any depth. Such low correlation coefficients can partially be attributed to modelling artefacts; however, they also suggest that not all of the  $V_s$  and  $Q_s$  anomalies in the continental upper mantle can be explained by  $T$  variations.

Global maps show that, by the sign of the anomaly,  $V_s$  and  $Q_s$  usually inversely correlate with lithospheric temperatures: most cratonic regions show high  $V_s$  and  $Q_s$  and low  $T$ , while most active regions have seismic and thermal anomalies of the opposite sign. The strongest inverse correlation is found at a depth of 100 km, where the attenuation model is best resolved. Significantly, at this depth, the contours of near-zero  $Q_s$  anomalies approximately correspond to the 1000 °C isotherm, in agreement with laboratory measurements that show a pronounced increase in seismic attenuation in upper mantle rocks at 1000–1100 °C. East–west profiles of  $V_s$ ,  $Q_s$  and  $T$  where continental data coverage is best (50°N latitude for North America and 60°N latitude for Eurasia) further demonstrate that temperature plays a dominant, but non-unique, role in determining the value of lithospheric  $V_s$  and  $Q_s$ .

At 100 km depth, where the resolution of seismic models is the highest, we compare observed seismic  $V_s$  and  $Q_s$  with theoretical  $V_s^T$  and  $Q_s^T$  values, respectively, that are calculated solely from temperature anomalies and constrained by experimental data on temperature dependencies of velocity and attenuation. This comparison shows that temperature variations alone are sufficient to explain seismic  $V_s$  and  $Q_s$  in *ca* 50 per cent of continental regions. We hypothesize that compositional anomalies resulting from Fe depletion can explain the misfit between seismic and theoretical  $V_s$  in cratonic lithosphere. In regions of active tectonics, temperature effects alone cannot explain seismic  $V_s$  and  $Q_s$  in the lithosphere. It is likely that partial melts and/or fluids may affect seismic parameters in these regions. This study demonstrates that lithospheric temperature plays the dominant role in controlling  $V_s$  and  $Q_s$  anomalies, but other physical parameters, such as compositional variations, fluids, partial melting and scattering, may also play a significant role in determining  $V_s$  and  $Q_s$  variations in the continental mantle.

**Key words:** composition, cratons, lithosphere, partial melts, Rayleigh waves, temperature.

1 INTRODUCTION

It is widely accepted that a significant part of lateral variations in seismic velocity and attenuation anomalies in the mantle can be attributed to temperature variations. Other possible causes of seismic anomalies include compositional variations, partial melts and/or fluids. In the present study we examine the correlations between seismic shear wave velocities ( $V_s$ ), inverse attenuation ( $Q_s$ ) and temperature ( $T$ ) in order to distinguish non-thermal and thermal origins of seismic anomalies in the continental uppermost mantle. Our special interest is to gain a better understanding of the physical nature of attenuation anomalies, because some recent interpretations of mantle attenuation assume that compositional variations in the mantle and/or the presence of partial melt have little effect on anelasticity and thus can be ignored (Nakajima & Hasegawa 2003). Numerous seismic and laboratory studies have addressed the question of the correlation between seismic elastic and anelastic properties, and temperatures.

1.1 The correlation between seismic velocity and attenuation

Anelasticity is manifested by the attenuation of the amplitude of waves during their propagation and is generally quantified in seismology by the quality factors,  $Q_p$  and  $Q_s$ , for compressional and shear waves, respectively. Global anelastic tomography studies (e.g. Romanowicz & Durek 2000) as well as regional seismic studies for the oceans (e.g. Sheehan & Solomon 1992) found that anomalies of seismic velocity and attenuation are correlated. Seismic studies of the Japan arc (Tsumura *et al.* 2000) demonstrated that low- $Q$  zones in the upper mantle are well correlated with low-velocity zones beneath active volcanoes where partial melting is expected. A strong inverse correlation of tomographic elastic and anelastic models beneath 100 km depth in the Tonga–Fiji region was used to derive an exponential empirical relationship between two of the parameters and to constrain the physical state of the upper mantle (Roth *et al.* 2000). A similarity between the observed relationship and predictions based on experimental studies (e.g. Sato *et al.* 1989; Karato 1993), permitted the authors to conclude that both attenuation and velocity anomalies should be primarily of thermal origin (Roth *et al.* 2000).

1.2 The correlation between seismic velocity and temperature

The strong effect of temperature on seismic velocity and elastic moduli has long been known from laboratory studies (e.g. Ide 1937; Birch 1943; Hughes & Cross 1951), implying that velocity perturbations in the uppermost mantle can be partly the result of temperature variations. Most of measurements were made at relatively low temperatures (< 600–700 °C) (Fielitz 1971; Kern 1978; Christensen 1979) with very limited data available for  $T > 1000$  °C (Berckhemer *et al.* 1982; Murase & Kushiro 1979; Murase & Fukuyama 1980; Sato *et al.* 1989; Jackson 1993). At high homologous temperatures  $T_m/T$  (where  $T_m$  is solidus temperature), a dependence of velocity on temperature is approximately linear with  $\partial V_s/\partial T = 0.35$  m s<sup>-1</sup> K<sup>-1</sup> (Sumino & Anderson 1982); however, at near-solidus temperatures both  $V_p$  and  $V_s$  seismic velocities rapidly decrease with an increase of a percentage of melt (Fig. 1; Murase & Fukuyama 1980).

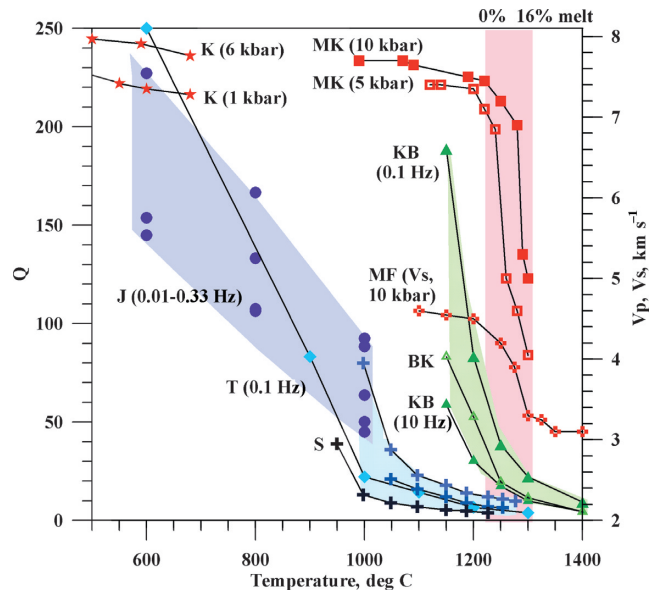


Figure 1. Experimental data on temperature dependence of  $Q$ ,  $V_p$  and  $V_s$ . Purple dots and purple shading (J):  $Q$  in dunite at  $P = 300$  MPa,  $f = 0.01$ – $0.33$  Hz (Jackson *et al.* 1992). Blue crosses and shading (S):  $Q$  in dry peridotite at  $P = 0.20, 0.48,$  and  $0.73$  GPa measured at ultrasonic frequencies (Sato & Sacks 1989). Green triangles and shading (BK, KB):  $Q$  in dunite at  $f = 0.1$  and  $10$  Hz (Berckhemer *et al.* 1982; Kampfmann & Berckhemer 1985). Light blue rhombs (T):  $Q$  in olivine aggregate at  $f = 0.09$  Hz (Tan *et al.* 1997). Red symbols: seismic velocities in peridotite. Stars (K):  $V_p$  at  $P = 1$  and  $6$  kbar (Kern 1978). Rectangles (MK):  $V_p$  at  $P = 5$  and  $10$  kbar (Murase & Kushiro 1979). Crosses (MF):  $V_s$  at  $P = 10$  kbar (Murase & Fukuyama 1980). Pink shaded area: the region where volume fraction of melt in peridotite increases from 0 to 16 per cent at  $P = 10$  kbar (Murase & Fukuyama 1980). At high temperatures, the effect of  $T$  on  $Q_s$  and  $V_s$  is different: a sharp increase in  $Q_s$  occurs at a lower temperature than a drop in  $V_s$  associated with partial melting.

High sensitivity of seismic velocities to temperature variations provides important constraints on the thermal regime of the upper mantle and has been successfully applied to estimate temperature variations in the upper mantle of Europe (Furlong *et al.* 1995; Sobolev *et al.* 1996; Goes *et al.* 2000) and North America (Goes & van der Lee 2002) from regional seismic tomography models. However, conversions of observed seismic velocities to temperatures should be treated with caution because non-thermal effects (resulting from compositional variations or the presence of melts/fluids) on seismic velocity variations are difficult to quantify (see Sections 3 and 5 below).

1.3 The correlation between seismic attenuation and temperature

The strong effect of temperature on seismic attenuation has been measured in a number of experimental studies (Fig. 1). Laboratory measurements of seismic wave attenuation in ultramafic upper-mantle rocks were carried out for a wide range of pressures and temperatures at ultrasonic frequencies (60–900 KHz) on dry peridotite (Sato *et al.* 1989), and at seismic frequencies (0.01–1 Hz) on dry dunite (Berckhemer *et al.* 1982; Jackson *et al.* 1992) and on synthetic polycrystalline olivine (Tan *et al.* 1997; Gribb & Cooper 1998; Jackson 2000; Jackson *et al.* 2002). They show that at seismic frequencies attenuation,  $Q^{-1}$ , in mantle rocks at subsolidus temperatures follows the Arrhenius law and exponentially increases with

temperature  $T$ :

$$Q^{-1} = A\tau^\alpha \exp(-\alpha E^*/RT), \quad (1)$$

where  $E^*$  is the activation energy,  $R$  is the gas constant,  $\tau$  is the oscillation period,  $A$  is a scaling constant and the exponent  $\alpha$  is approximately 0.15–0.30 as determined from seismic studies and laboratory measurements on the upper-mantle rocks (e.g. Jackson *et al.* 1992).

Laboratory measurements of the temperature dependence of seismic attenuation for olivine-rich rocks were used in a number of studies where mantle geotherms were calculated from seismic attenuation (e.g. Kampfmann & Berckhemer 1985; Sato *et al.* 1988; Sato & Sacks 1989; Karato 1990; Nakajima & Hasegawa 2003). However, attenuation values determined from laboratory studies (eq. 1) include only intrinsic (anelastic) attenuation, while global and regional seismological studies measure the sum of anelastic and scattering effects.

The high sensitivity of seismic attenuation to temperatures, revealed in laboratory measurements, was recently used to correlate seismic attenuation anomalies in the mantle with temperature variations (e.g. Durek *et al.* 1993; Romanowicz 1994; Mitchell 1995; Romanowicz 1995). Regional seismic studies of North Atlantic (Sheehan & Solomon 1992) and western Pacific subduction zones (Nakanishi 1978; Bodri *et al.* 1991; Tsumura *et al.* 2000) suggest that a significant part of attenuation anomalies in the oceanic upper mantle can be explained by temperature variations and partial melting. In a global study, Romanowicz (1995) found that in the depth interval from 100 to 300 km shear wave attenuation  $Q_s^{-1}$  calculated from the amplitudes of low-frequency Rayleigh waves on a  $10^\circ \times 10^\circ$  grid (the QR19 model) correlates with the surface heat flow values (the correlation coefficient  $r$  is in the range 0.20 to 0.35 with the correlation peak at  $z \sim 200$  km); while at depths of  $z > 250$  km  $Q_s^{-1}$  correlates with the hot spot distribution, with the strongest correlation ( $r \sim 0.35$ – $0.38$ ) at  $z \sim 300$ – $660$  km. However, it is important to note that these correlations are based mostly on high attenuation anomalies documented for oceanic regions (mid-ocean ridges, the western Pacific subduction zones and the East Pacific Rise), but not for the continents.

On the continents, studies of seismic shear wave attenuation revealed a qualitative correlation between regional attenuation anomalies in the crust and upper mantle and tectonic provinces (e.g. Nakanishi 1978; Canas & Mitchell 1978; Roullet 1982; Dziewonski & Steim 1983; Chan & Der 1988; Mitchell *et al.* 1997; Sarker & Abers 1998; Selby & Woodhouse 2002). Low seismic attenuation in the upper mantle has been reported beneath western Australia, western Africa and the Himalayas in a number of recent anelastic tomographic models (e.g. Romanowicz 1994, 1995; Bhattacharyya *et al.* 1996), supporting the idea of the thermal origin of most of the upper-mantle  $Q_s^{-1}$  anomalies.

Few studies have been carried out so far to analyse the correlations between all three parameters ( $V_s$ ,  $Q_s$  and  $T$ ) on the global scale. Here, we aim to evaluate the influence of  $T$  on  $V_s$  and  $Q_s$  for the continents and compare shear wave velocities and attenuation at different depths in the upper 150 km of the mantle with temperature estimates (Artemieva & Mooney 2001). The present work uses a new 3-D model of  $Q_s$  (Billien *et al.* 2000) with the maximum resolution at a depth of 100 km and the corresponding 3-D model of  $V_s$  (both  $Q_s$  and  $V_s$  are effectively degree-13 spherical harmonic; Billien 1999).

We first calculate global correlation coefficients between all pairs of the parameters and find that, though a large part of seismic anomalies correlates with  $T$  variations in the upper mantle, temperatures

alone cannot explain the amplitudes of velocity and attenuation anomalies. We discuss other mechanisms responsible for  $V_s$  and  $Q_s$  variations and examine a qualitative correlation between  $V_s$ ,  $Q_s$  and  $T$  maps for 50, 100 and 150 km depths for different tectonic provinces. We finally use experimental relationships between seismic parameters and temperature and calculate global maps of mantle attenuation and shear velocity predicted from temperatures. A comparison of these maps with the seismic maps permits the outlining of regions where mechanisms other than thermal should be involved to explain seismic anomalies.

## 2 METHODS FOR CALCULATING $V_s$ , $Q_s$ AND $T$ IN THE UPPER MANTLE

### 2.1 Shear wave velocity and attenuation models

#### 2.1.1 The data

Studies of seismic attenuation in the Earth's mantle using surface waves began with the works of Anderson & Archambeau (1964), Ben-Menahem (1965) and Anderson *et al.* (1965): however, global attenuation models have been calculated only recently (Romanowicz 1995; Bhattacharyya *et al.* 1996; Durek & Ekström 1996; Reid *et al.* 2001; Selby & Woodhouse 2002). Velocity and attenuation maps used in this study are based on surface waves that sample the Earth all along their path from the surface to deep upper mantle. Compared to body waves, which have poor sampling (especially in the oceans) and thus very poor vertical resolution in the upper mantle, surface waves provide good vertical resolution down to 200–300 km and have acceptable horizontal resolution for discussing large geological units.

In the present study, we use a subset of the 3-D upper-mantle  $V_s$  and  $Q_s$  models of Billien (1999) and Billien *et al.* (2000). The models are based on the inversion of broad-band seismological data of fundamental mode surface waves. The data used are amplitude and frequency spectra at periods between 40 and 150 s for the direct path Rayleigh wave data (the first orbit only), which is a subset of the extensive phase and amplitude data set measured and selected by Trampert & Woodhouse (2001). For a detailed description of the automatic measurement technique and implementation of the strict data rejection criteria, see Trampert & Woodhouse (1995).

#### 2.1.2 Simultaneous inversion for shear wave velocity and attenuation

$V_s$  and  $Q_s$  calculations are based on a two-step inversion (for details see Billien 1999; Billien *et al.* 2000). First, from phase ( $\Phi$ ) and amplitude ( $A$ ) measurements, Billien *et al.* (2000) made a simultaneous regionalization of the phase velocity ( $c_R$ ) and the surface wave attenuation factor ( $q_R = Q_R^{-1}$ ) at several periods:

$$A_i(\omega), \Phi_i(\omega) \Rightarrow c_R(\omega, \theta, \phi), q_R(\omega, \theta, \phi). \quad (2)$$

A depth inversion of the phase velocities and surface wave attenuation maps at different frequencies provides a 3-D model of the shear wave velocity  $V_s$  and the shear wave attenuation factor  $q_s$  ( $q_s = Q_s^{-1}$ ) (Billien 1999):

$$q_R(\omega, \theta, \phi) \Rightarrow q_s(z, \theta, \phi), \quad (3)$$

$$c_R(\omega, \theta, \phi) \Rightarrow V_s(z, \theta, \phi). \quad (4)$$

The classical inversion method of Tarantola & Valette (1982) is used at this stage. It employs a least-squares method and includes information on the expected model in the form of: (i) an *a priori* model and (ii) an *a priori* covariance of the model. Billien (1999) used a layered version of PREM with 50 km thick layers in the depth range 25–825 km as the *a priori* velocity and attenuation models. The upper 25 km (the crustal layer) are fixed to the PREM value and are never updated in the inversion process. In the globally spherical PREM model (Dziewonski & Anderson 1981) the radial average  $Q_s$  has a value of 600 at depths between 40 and 80 km, and  $Q_s = 80$  between 80 and 220 km. Similar to the attenuation model, the velocity value in the crustal layer is not allowed to change in the inversion, because the phase velocity maps are corrected at all periods for the crustal effects calculated from the CRUST5.1 model (Mooney *et al.* 1998). These corrected maps, as a function of period, are then inverted to obtain the 3-D  $S$ -velocity model.

The *a priori* covariance of the attenuation model is described by *a priori* standard deviations as a function of depth. In order to allow larger  $Q_s$  variation in the lithosphere, which is believed to be more heterogeneous than the asthenosphere, a standard deviation on the *a priori* model  $\sigma_m$  was chosen to be twice as large in the lithosphere as in the asthenosphere (Billien 1999).

### 2.1.3 Resolution of the $V_s$ and $Q_s$ models

The problem of the effects of focusing/defocusing (as a result of refraction at strong velocity heterogeneities) on the wave amplitude was resolved by Billien (1999) by performing simultaneously a global inversion of the two seismic parameters,  $V_s$  and  $Q_s$ : it ensured full compatibility between the final  $V_s$  and  $Q_s$  models in terms of focusing/defocusing effects. Specifically, a linear approximation (Woodhouse & Wong 1986) was used to include the focusing effect in the inversion.

The observed attenuation is a mixture of scattering (largely as a result of reflection at velocity heterogeneities) and intrinsic (anelastic) attenuation. Because both of them lead to an exponential decay of amplitude with time (or epicentral distance; Billien 1999), they cannot be separated by inversion and a contribution of scattering into the observed attenuation cannot be completely ruled out even in the frequency range used in this study (Billien 1999).

In the regionalization, smooth models were favoured over rough models. This is achieved by use of a cost function based on the norm of the Lagrangian (second derivatives) of the model (Billien *et al.* 2000) rather than a classical damped least-squares method based on the norm of the model (e.g. Tarantola & Valette 1982). In our model, phase velocity and attenuation maps, and then  $V_s$  and  $Q_s$  maps, are represented by spherical harmonics up to degree  $l = 20$ . However, the effective degree of the computed maps is only 12 to 13, as a result of the damping used in the regionalization that progressively shuts off the highest degree terms. This effective degree corresponds to an *ca* 3000 km wavelength of attenuation and velocity lateral variations.

The vertical resolution of  $V_s$  and  $Q_s$  models is *ca* 50 km. The depth resolution matrix shows that the best resolved layer for  $Q_s$  is at a depth of  $100 \pm 25$  km. Because Billien (1999) and Billien *et al.* (2000) use a damped inversion, the amplitude of  $Q_s$  variations at each depth layer within the lithosphere is probably underestimated. The tuning of the cost function in the first step of the inversion largely influences the absolute values of the  $Q_s$  variations (Billien *et al.* 2000) and thus the values of  $Q_s$ , but the sign of the variations and the sharpness of the contrasts are less sensitive to this tuning.

Both attenuation and velocity anomalies are calculated with respect to PREM (Figs 2–4), which was used as a reference model

in the seismic inversion, because both  $V_s$  and  $Q_s$  models are global. Usage of PREM does not have a strong effect on calculated  $Q_s$  because focusing/defocusing effects depend not on absolute velocity values, but on the velocity gradient (Billien 1999). If a continent-dominated reference model were used in the seismic inversion, it would have led to a systematic shift of the absolute values, but would not have affected the general pattern of the anomalies. The difference in  $V_s$  between PREM and a preferable continent-dominated model is *ca* 2–2.5 per cent at the depth 100–150 km (Ritsema, private communication, 2003). The ocean-dominated PREM model led to a systematic shift to higher  $V_s$  on the continents, as can be clearly seen in Figs 2(a), 3(a) and 4(a), where white near-zero-anomaly regions correspond, in fact, to 0 to +3 per cent velocity anomalies. Because we are primarily interested in the patterns of velocity and attenuation variations and their correlations with temperature anomalies, the baseline of the reference model is not important for the further analysis.

## 2.2 Calculation of lithospheric temperatures

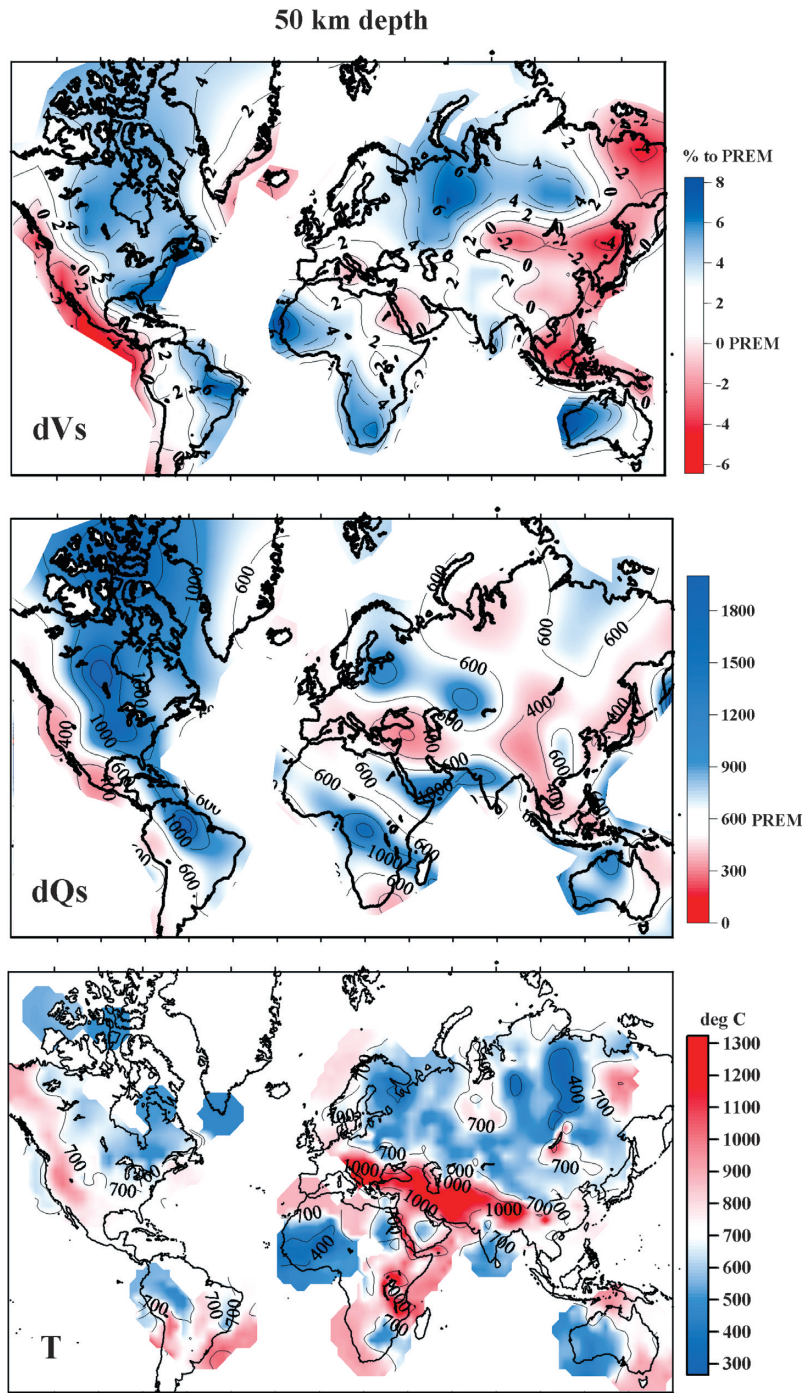
### 2.2.1 The model

The thermal state of the continental lithosphere has been addressed in numerous publications (see Artemieva & Mooney 2001, for references). However, most of these are regional studies and a lack of consistent assumptions on crustal heat production and thermal conductivity makes it impossible to combine them into a global model.

Two global studies of the thermal regime of the continental lithosphere are available at present. The model of Pollack & Chapman (1977) is based on a degree-12 spherical harmonic representation of global surface heat flow and provides an estimate of typical continental and oceanic geotherms and a rather generalized global map of lithospheric thermal thickness. In that model, continental geotherms were calculated assuming that 40 per cent of the mean surface heat flow arises from a near-surface, enriched radiogenic zone with a uniform 8 km thickness. Subsequent regional thermal and xenolith studies have proved that, on the whole, the set of geotherms calculated by Pollack & Chapman (1977) provides a good approximation of lithospheric temperatures. However, in some cratonic regions the assumptions behind this global model are invalid (e.g. Jaupart & Mareschal 1999).

Because cratons are expected to produce the strongest negative temperature anomalies within the continental lithosphere, here we use a more recent global thermal model for stable continents (Artemieva & Mooney 2001), in which regional seismic data on the crustal structure together with regional geological, petrological and laboratory data were used to constrain the vertical distribution of crustal heat production in different terranes (with sizes varying from approximately  $1^\circ \times 1^\circ$  to  $5^\circ \times 5^\circ$ ) and its total contribution to the surface heat flow. Models of the lateral and vertical distribution of heat production and thermal conductivity in the continental lithosphere used to constrain the thermal model are discussed in detail in Artemieva & Mooney (2001) and we refer the reader to the original publication for a detailed description of the data, the method and the calculation scheme. The consistency of assumptions for the depth distributions of thermal parameters in the thermal model permits the quantitative comparison of lithospheric geotherms, even if the calculated values contain systematic biases.

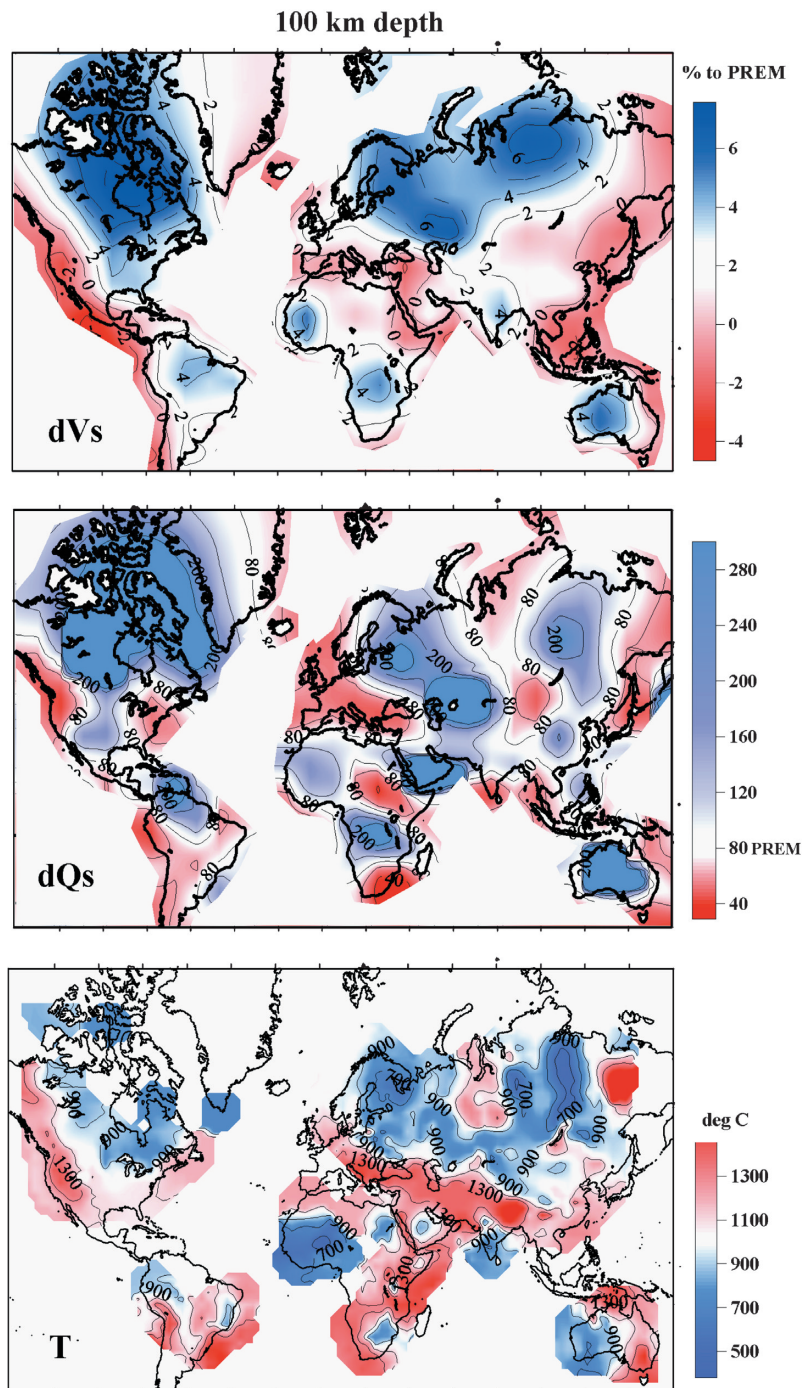
Estimates of the thermal regime of the stable continental lithosphere are based on the solution of the steady-state thermal conductivity equation, which allows the calculation of temperatures  $T(z)$  within the crust and lithospheric mantle, constrained by borehole



**Figure 2.** Map of shear wave velocity  $V_s$  (as variation in per cent relative to PREM), inverse attenuation  $Q_s$ , and temperature  $T$  (in  $^{\circ}\text{C}$ ) for continental lithosphere at a depth of  $50 \pm 25$  km. Note that for the Tibetan plateau and the Andes, where the crustal thickness exceeds 60–70 km, the anomalies refer to the crustal depth, rather than to the subcrustal lithosphere. Blue colors correspond to cold regions with high  $V_s$ , high  $Q_s$  and low  $T$ . Red colors correspond to hot lithosphere with low  $V_s$ , low  $Q_s$  and high  $T$ . The dominant pattern on all three maps is the correlation of blue colors with stable Precambrian cratons and red colors with tectonically active Phanerozoic regions. The maps show a correlation in blue regions in the Canadian, Baltic, Siberian, West African and Australian cratons. Red regions that correlate include western North America, Europe, the Middle East, eastern China and eastern Australia.

heat flow measurements and the distribution of thermal parameters with depth. The global heat flow database (Pollack *et al.* 1993) updated for heat flow measurements reported since 1993 formed the basis for calculations of temperatures in the continental upper mantle. The lack of complete information on vertical and lateral distri-

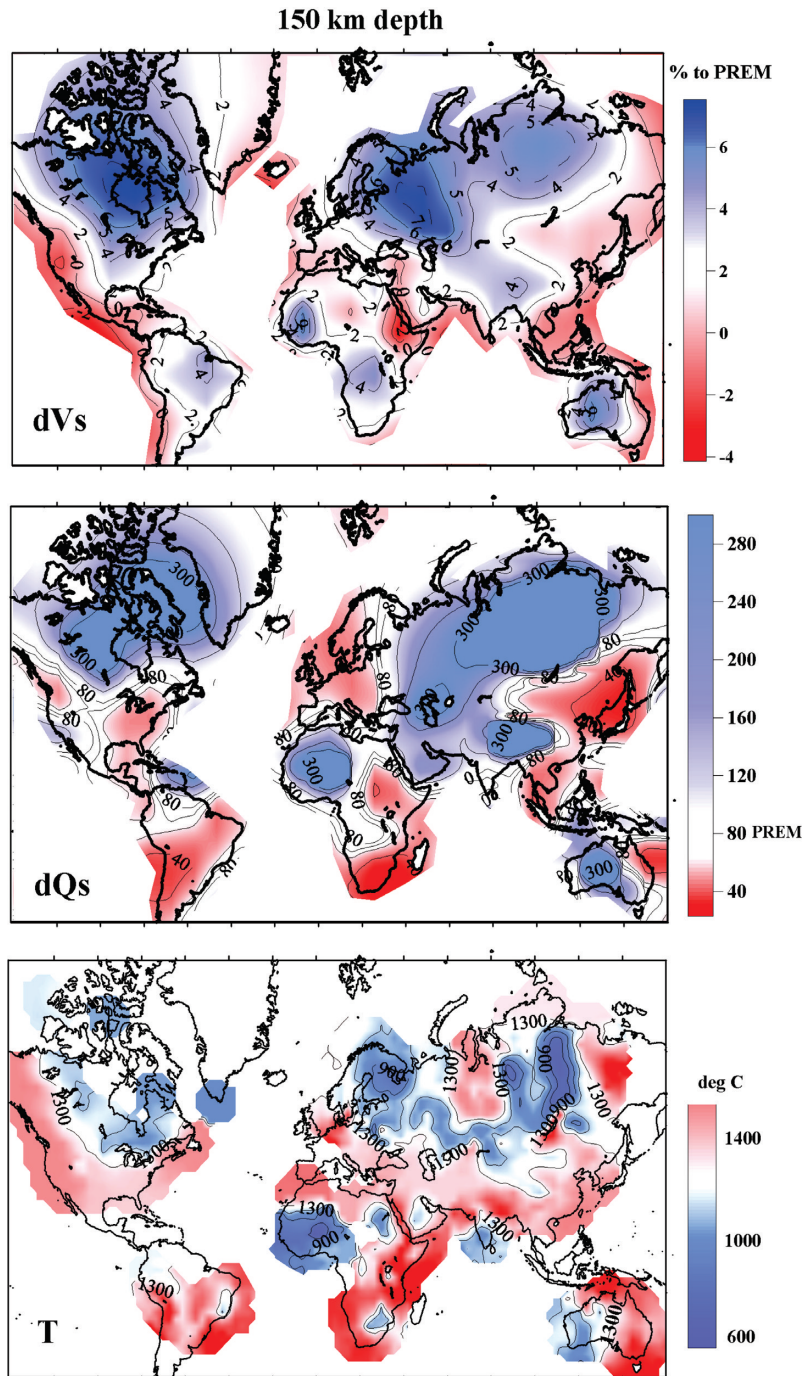
bution of crustal heat production (which is one of the key parameters controlling the resultant thermal models of the lithosphere), as well as the very uneven distribution of heat flow borehole data do not justify a 2-D or 3-D thermal modelling on a global scale. Errors associated with high lateral heat conduction can be large only in



**Figure 3.** Map of shear wave velocity  $V_s$  (as variation in per cent relative to PREM), inverse attenuation  $Q_s$  and temperature  $T$  (in  $^{\circ}\text{C}$ ) for continental lithosphere at a depth of  $100 \pm 25$  km. Presentation as in Fig. 2. There is a close correspondence in the anomaly patterns for most of the Precambrian cratons and the Cenozoic regions, especially for the Precambrian Eurasia, West Africa, northern Canada, western North America and central Australia. This correspondence indicates that temperature plays an important role in determining the strength of  $V_s$  and  $Q_s$  anomalies at this depth in the lithosphere.

regions with a strong contrast in heat production and/or conductivity (Jaupart 1983). However, assigning unconstrained values of thermal parameters and surface heat flow in areas with scarce or non-existing data would have introduced uncontrolled errors into the entire model.

The estimates of the lithospheric thermal thickness are in a good agreement with global and regional seismic tomography results, and with regional xenolith geotherms (*cf.* Artemieva & Mooney 2001). The maps of temperature distribution in the continental lithosphere at the depths of 50, 100 and 150 km form the basis of our comparison



**Figure 4.** Map of shear wave velocity  $V_s$  (as variation in per cent relative to PREM), inverse attenuation  $Q_s$  and temperature  $T$  (in  $^{\circ}\text{C}$ ) for continental lithosphere at a depth of  $150 \pm 25$  km. Presentation as in Fig. 2. All low temperatures (lower panel) correspond to stable cratons. In contrast, high  $V_s$  and high  $Q_s$  anomalies include these same stable cratons plus some regions of active tectonics, such as the southern Eurasia suture zone extending from the Alps to the Tibetan plateau. This may indicate the presence of cold subducted lithosphere beneath the southern European suture zone. High mantle temperatures ( $>1300$   $^{\circ}\text{C}$ , shown by red colors in Fig. 4c) dominate the non-cratonic parts of the continents and imply a lithospheric thickness less than 150 km, if defined as the conductive layer above the mantle adiabat of 1300  $^{\circ}\text{C}$ .

of lateral variations of temperature,  $V_s$  and  $Q_s$  at different depths in the continental lithosphere.

**2.2.2 Resolution of the  $T$  models**

The constraints provided by the thermal model are reliable for stable continental regions, where the assumption of a steady-state thermal regime of the lithosphere is valid. For most tectonically active re-

gions (e.g. the Alpine–Himalayas region, the Andes and western USA) the temperature distribution in the upper mantle is not calculated from surface heat flow but is based on available petrologic and non-steady-state constraints (e.g. Lachenbruch & Sass 1977; Polyakov *et al.* 1988; Henry *et al.* 1997; Le Pichon *et al.* 1997). Lithospheric temperature maps are calculated from point (i.e. borehole) data (for heat flow data coverage see Fig. 5) and give

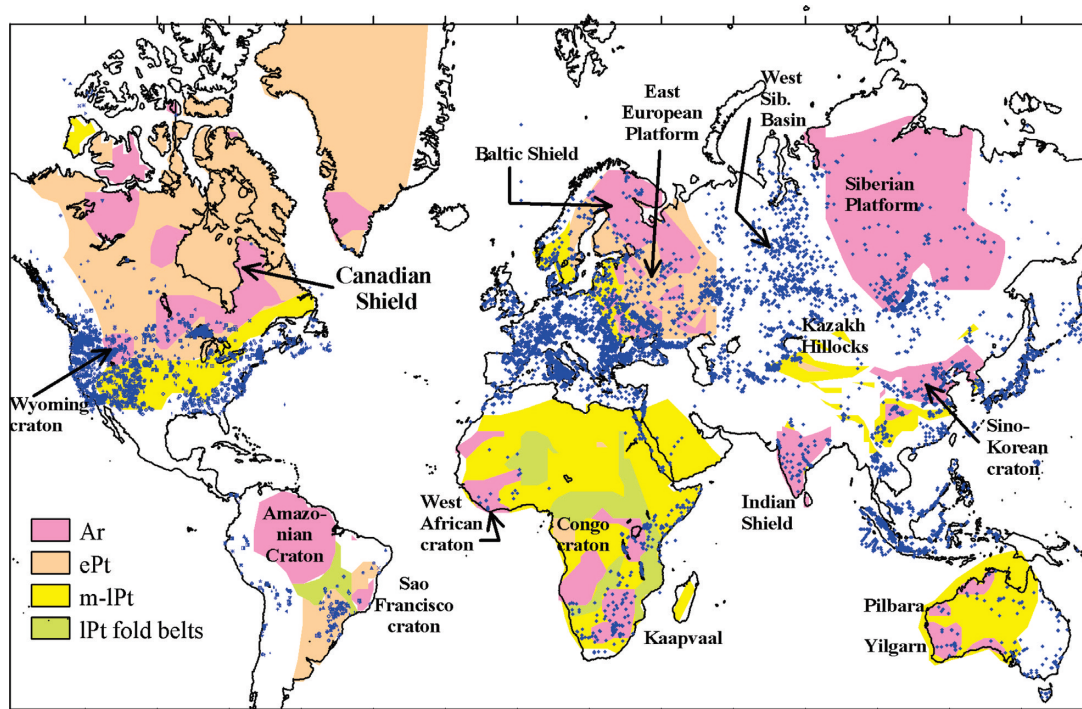


Figure 5. Location of Precambrian cratons and sites of borehole heat flow measurements on the continents [based on global heat flow compilation (Pollack *et al.* 1993) updated for heat flow measurements reported in 1993–1999].

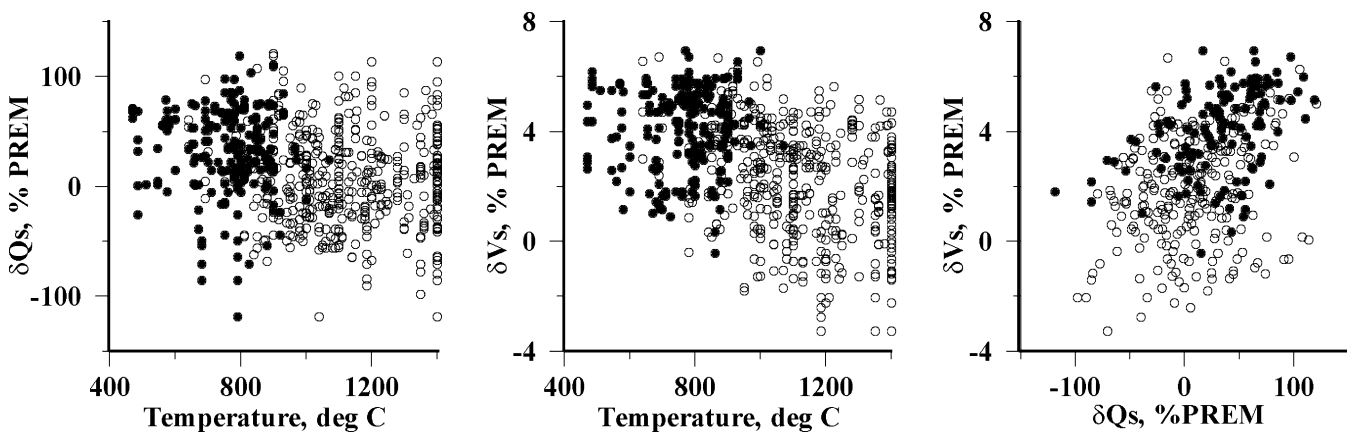


Figure 6. Correlations between shear wave velocity  $V_s$  (as variation in per cent relative to PREM), inverse attenuation  $Q_s$  (as variation in per cent relative to PREM) and temperature  $T$  (in  $^{\circ}\text{C}$ ) for continental lithosphere at a depth of 100 km.  $V_s$  and  $Q_s$  values are for the sites of thermal modelling. All data for the continents (open circles); data for Archean and early Proterozoic cratons only (solid circles). The calculated correlation coefficients are less than 0.42 for any pair of the parameters (see Table 1). The correlation is especially bad for the Archean–early-Proterozoic cratons, suggesting that a significant part of lateral  $V_s$  and  $Q_s$  variations in cratonic lithosphere has a non-thermal origin.

estimates of  $T(z)$  in the lower lithosphere with an accuracy of *ca*  $100^{\circ}\text{C}$  (see Artemieva & Mooney 2001, for sensitivity estimates). These local geotherms calculated for regions with typical sizes of  $1^{\circ} \times 1^{\circ}$  to  $5^{\circ} \times 5^{\circ}$  were interpolated by krigging with an interpolation radius of  $10^{\circ}$  (Figs 2c, 3c and 4c) that preserved the accuracy of  $T(z)$  estimates in the vicinity of heat flow boreholes and provided smoothed anomaly contours between them.

Note that interpolation of local geotherms introduces artefacts in the form of unconstrained haloes in regions with scarce borehole data (e.g. the Arabian shield, Central Africa, South America,

northern Canada and Greenland) that can extend as far as 1000 km from sites of heat flow measurements. In regions with a strong contrast in tectonic and thermal regimes (e.g. the transition between stable continental USA and western North America, or the transition between Precambrian and Variscan Europe, or East African rift system), an interpolation resulted in a smeared thermal contrast. Thus, the spatial resolution of the temperature maps is uneven, but is within a range of 200–1000 km and is significantly higher than the lateral resolution of seismic attenuation and shear wave velocity maps.

**Table 1.** Correlation coefficients between seismic velocity, attenuation, temperature and surface heat flow at different depths in the continental upper mantle.

Parameters	$z = 50$ km	$z = 100$ km	$z = 150$ km
$Q_s$ vs. $T$ (*)	-0.22	-0.33	-0.47
$\ln Q_s$ vs. $1/T$ (*)	0.27	0.31	0.36
$V_s$ vs. $T$ (*)	-0.31	-0.56	-0.56
$V_s$ vs. $Q_s$ (*)	0.38	0.66	0.41
$Q_s$ vs. surface heat flow (*)	-0.19	-0.18	-0.04
$V_s$ vs. surface heat flow (*)	-0.26	-0.39	-0.36
$Q_s$ vs. $T$ (black and white)	-0.24	-0.24	-0.25
$V_s$ vs. $T$ (black and white)	-0.20	-0.36	-0.26
$V_s$ vs. $Q_s$ (black and white)	0.23	0.48	0.28
$Q_s$ vs. $T$ (all data)	-0.19	-0.23	-0.34
$V_s$ vs. $T$ (all data)	-0.19	-0.38	-0.42
$V_s$ vs. $Q_s$ (all data)	0.38	0.41	0.41

(\*) At each depth, values within the band around the average value (zero anomaly), approximately corresponding to white areas in Figs 2–4, are excluded from the analysis. The width of the band is 10 per cent of the total range of parameter variations. The correlations for  $V_s - Q_s$  and  $Q_s - T$  at  $z = 50$  and 150 km and for  $V_s - T$  at  $z = 50$  km should be treated with caution as a result of low resolution of  $V_s$  and  $Q_s$  models at these depths. For a comparison, we calculated the correlation between the signs of the anomalies only (black and white models) and excluding the 10-per-cent-wide band around the average value from the statistical analysis. Lower correlations indicate that the amplitudes of the anomalies are important in global correlations. Correlation coefficients calculated when all data were included into the analysis are presented for a comparison as well.

### 3 QUANTITATIVE CORRELATIONS

#### $V_s - Q_s - T$

##### 3.1 Calculation of correlation coefficients

Laboratory evidence for a strong correlation between seismic parameters and temperatures has recently been used in a number of studies to constrain mantle temperatures from elastic and anelastic tomography (e.g. Sobolev *et al.* 1996; Sato *et al.* 1988). We next discuss a validity of these approaches and examine the quantitative correlation between each pair of the parameters  $V_s$ ,  $Q_s$  and  $T$  at depths 50, 100 and 150 km (Table 1). To calculate the correlation coefficients ( $r$ ), seismic data were converted from spherical harmonic representation to a  $1^\circ \times 1^\circ$  spatial grid. Two interpolations of temperatures were calculated from local geotherms: (i) with a  $5^\circ$  radius to reduce artefacts of interpolation in regions with sparse data and (ii) with a  $30^\circ$  radius to bring the temperature model to a resolution comparable to the resolution of seismic models. These interpolated temperature maps were resampled to a  $1^\circ \times 1^\circ$  spatial grid. In the case of a  $30^\circ$  interpolation, correlation coefficients between seismic parameters and  $T$  are *ca* 0.05 weaker than in case of a  $5^\circ$  interpolation, probably as a result of interpolation haloes unconstrained by heat flow data: thus, we further discuss only the results for a  $5^\circ$  interpolation of mantle geotherms.

Correlations between the parameters  $V_s$ ,  $Q_s$  and  $T$  were calculated on a  $1^\circ \times 1^\circ$  grid: for each of the parameters the values were normalized by the maximum amplitude calculated for each depth. We first included all  $V_s$ ,  $Q_s$  and  $T$  data into the analysis. The correlation coefficients were very low, less than 0.42 for any pair of the parameters and at any depth and, when plotted as  $x$ - $y$  plots, formed clouds (Fig. 6). Roth *et al.* (2000) found similar clouds when comparing global attenuation and velocity models [QR19 of Romanowicz (1995) and S12.WM13 of Su *et al.* (1994)]. The correlation is in particular low for the Archean-early Proterozoic cratons, suggesting

that for these regions a significant part of lateral  $V_s$  and  $Q_s$  variations has a non-thermal origin. Because the correlation coefficient is sensitive to the sign of seismic and temperature anomalies when their amplitudes are small, we next excluded from the correlation analysis the bands around the zero-value anomalies. The width of the bands was chosen as 10 per cent of the total range of parameter variations around the average value (zero anomaly): it approximately corresponds to the white areas in Figs 2–4 and is close to uncertainties in amplitudes in seismic and temperature models.

##### 3.2 $V_s - T$ correlation

Strong effect of temperature on seismic velocities is well-known from laboratory studies (Fig. 1). However, our statistical analysis of the  $V_s - T$  correlation in the upper mantle of the continents suggests that even though there is the expected negative correlation between the two parameters ( $r = -0.56$  at  $z = 100$  and 150 km; Table 1), the conversion of seismic velocities into temperatures (e.g. Sobolev *et al.* 1996; Goes *et al.* 2000; Goes & van der Lee 2002) is subject to considerable uncertainty. Our data show that at a depth of 100 km, the range of possible temperatures corresponding to any given velocity value is at least  $\pm 250$  °C (Fig. 6). This indicates that a significant part of the velocity anomalies in the continental upper mantle can be the result of physical mechanisms other than thermal (i.e. compositional heterogeneities, anisotropy, fluids, or melts). With the present resolution of seismic and thermal models it is, however, difficult to say which part of the  $V_s$  variations has a true non-thermal origin. We further address this question in Section 5.3. Here we note that the maximum values of  $V_s$  anomalies are found for the Canadian shield and the East European and Siberian cratons. Although recent studies suggest that the temperature effect on velocities is much stronger than iron content variations (Deschamps *et al.* 2002), compositional difference between the Precambrian and Phanerozoic subcrustal lithosphere (e.g. Boyd 1989; Griffin *et al.* 1999) can be partly responsible for low  $V_s - T$  global correlation on the continents.

##### 3.3 $Q_s - T$ correlation

An exponential law for the correlation of temperature and seismic attenuation has long been known from laboratory experiments (Fig. 1). Sato *et al.* (1988) were among the first to apply it to constrain oceanic geotherms from seismic attenuation and laboratory data. Other attempts to calculate temperatures from seismic attenuation data include studies by Sarker & Abers (1999) for the crust of south-central Eurasia and by Nakajima & Hasegawa (2003) for the mantle wedge beneath northeastern Japan. Though seismic data show the existence of a regional correlation between deep temperatures and attenuation (e.g. Bodri *et al.* 1991), the statistical analysis (Table 1) shows that, globally, the correlation coefficient peaks at  $r = -0.47$  at  $z = 150$  km and is much weaker at shallower depths (e.g.  $r \sim -0.33$  at  $z = 100$  km, where the resolution of the attenuation model is the best). Because experimental studies suggest a logarithmic dependence of  $Q_s$  on  $T$  (eq. 1), we have also calculated the correlation coefficient between  $\ln(Q_s)$  and  $1/T$  and found it to be very close to the  $Q_s - T$  correlation. A poor global correlation between  $Q_s$  and  $T$  on the continents (less than 0.47 at any depth, regardless of the way it has been calculated) suggests that, as for  $V_s$ , effects other than temperature (e.g. composition, fluids, partial melt) play an important role in producing attenuation anomalies in the upper mantle.

In stable continental regions, where a presence of partial melts or fluids at a 100 km depth is unlikely, one would expect that

compositional variations could play the dominant role in smearing the correlation between  $Q_s$  and  $T$ . However, no experimental studies have examined how depletion can influence seismic attenuation. Recent laboratory measurements of attenuation in MgO (Getting *et al.* 1997) suggest that a difference in Mg content in the subcrustal lithosphere of the cratons and Palaeozoic platforms can be reflected in  $Q_s$  values. Because cratonic lithosphere was formed by a collage of continental and oceanic terranes with different structure and bulk composition, large-scale lithospheric heterogeneities can produce scattering and also be responsible for a lack of  $Q_s$ - $T$  correlation. A regional study of the Japan islands has demonstrated that scattering can account for 10–30 per cent of the measured attenuation (Hoshiya 1993). Nevertheless, if variations of scattering in the upper mantle are small, the correlation between intrinsic (and total) attenuation and temperature will still hold.

### 3.4 $Q_s$ - $V_s$ correlation

Because laboratory measurements on olivine-rich rocks suggest that both velocity and attenuation are strongly temperature dependent, one may expect  $V_s$  and  $Q_s$  to be correlated, as indicated by regional (e.g. Sheehan & Solomon 1992; Tsumura *et al.* 2000) and global studies (Romanowicz 1995). For example, for the Fiji–Tonga region, Roth *et al.* (2000) found that, at depths greater than 100 km, both anomalies are strongly correlated and can be fit by an exponential equation; at shallower depths their data show a high degree of scatter. They also examined the correlation between  $V_s$  and  $Q_s$  anomalies derived from global seismic tomography studies (Su *et al.* 1994; Romanowicz 1995) and concluded that the strong correlation between the two parameters in the Tonga subduction zone is more a regional phenomenon than a general rule. The statistical analysis of a global  $Q_s$ - $V_s$  data set for the continents presented here (Table 1) supports this conclusion. At 100 km depth we observe the correlation  $r = 0.66$  between  $Q_s$  and  $V_s$  in the upper mantle of the continents: this is the highest correlation we found for any pair of the parameters. Higher sensitivity of  $Q_s$  than  $V_s$  to temperature variations (e.g. Berckhemer *et al.* 1982; Karato & Spetzler 1990; Jackson *et al.* 1992) can be partly responsible for the moderate global correlation between  $Q_s$  and  $V_s$ .

### 3.5 Why are correlation coefficients so low?

Low correlation coefficients between  $V_s$ ,  $Q_s$  and  $T$  determined in statistical analysis for the continents contradict laboratory data on a strong dependence of  $V_s$  and  $Q_s$  on  $T$ . However, our correlation coefficients for 100 km depth (where high resolution is expected for both  $V_s$  and  $Q_s$  models) are significantly higher ( $r = -0.56$  and  $-0.33$ , correspondingly) than found for a global correlation between the QR19 attenuation model and surface heat flow ( $r = 0.20$ – $0.35$ ) at  $z < 250$  km (Romanowicz 1995). On the other hand, when the oceanic upper mantle is excluded from the analysis, the correlation between  $Q_s$  model and surface heat flow becomes weaker ( $r \sim -0.20$ ) (Table 1).

Relatively weak global correlations between thermal and seismic anomalies in the continental upper mantle suggest that, globally, temperatures alone cannot explain a large part of attenuation and velocity variations. A set of different physical mechanisms as well as artefacts associated with inversion, resolution and interpolation can deteriorate the global  $V_s$ - $Q_s$ - $T$  correlations from what can be expected from experimental studies.

- (i) Accuracy of seismic and thermal models used in this study:
- (a) Attenuation measured in laboratory includes only anelastic at-

tenuation (eq. 1). However, the attenuation model used in the present study also includes scattering effects, which could not have been separated during the inversion.

(b) Uncertainties in  $V_s$  and  $Q_s$  models can further affect calculated correlation coefficients. In some regions, strong offshore seismic anomalies are smeared into onshore regions.

(c) Errors in temperature estimates can be high in tectonically active regions, where temperature constraints from data on regional metamorphism were extrapolated to mantle depths.

(d) Seismic models provide a snap-shot of the physical state of the upper mantle, while surface heat flow reflects its past thermal regime. In some regions this can produce apparently non-correlated seismic and thermal anomalies.

(e) Seismic and thermal models have different lateral resolution and different spatial coverage. Because  $Q_s$  and  $V_s$  models were calculated in simultaneous inversion from the same seismic data set and have the same lateral resolution, one would expect a strong  $Q_s$ - $V_s$  correlation, if  $Q_s$  and  $V_s$  were not affected by thermal, compositional and fluid regime variations in a different way. Thus, low correlation coefficients can hardly be attributed solely to resolution problems and uncertainties in seismic and thermal models and raise the possibility that non-thermal effects on  $V_s$  and  $Q_s$  are significantly larger than usually assumed.

(ii) Anisotropy: The effect of anisotropy on surface waves at the periods used in the present study was examined recently by Trampert & Woodhouse (2003). However, our seismic models do not account for azimuthal anisotropy in the upper mantle. This oversimplification may introduce biases in seismic models. Such biases are generally small and their effects on the velocity and attenuation contrasts are likely to be even smaller, and thus should not significantly affect  $V_s$ - $Q_s$ - $T$  correlations.

(iii) Variations in composition and mineralogy: The effect of mantle composition and mineralogy on  $V_s$  has been previously studied in detail (e.g. Nolet & Zielhuis 1994; Jackson & Rigden 1998; Goes *et al.* 2000; Deschamps *et al.* 2002). Based on forward calculations, these authors conclude that variations of iron content in the lithospheric mantle have a larger effect on  $V_s$  than any other variations in composition and mineralogy. However, their effect is significantly smaller than the effect of  $T$  variations: for example, a 1 per cent  $V_s$  anomaly can be explained either by a 4 per cent variation of iron content or by a thermal anomaly of 50–100 °C (Nolet & Zielhuis 1994; Deschamps *et al.* 2002). Little is known of the effect of MgO or FeO variations on the attenuation (Getting *et al.* 1997). However, the results of Billien *et al.* (2000) and Sarker & Abers (1998) show a systematic decrease of  $Q_s$  with tectonothermal age of continental lithosphere. Thus, we cannot rule out the possibility that compositional variations in continental upper mantle can affect  $Q_s$ - $T$  and  $V_s$ - $Q_s$  correlations.

(iv) Partial melts and fluids: Stable continental lithosphere is drier compared with oceanic lithosphere (Hirth *et al.* 2000). Thus, in >100 km thick lithosphere of Palaeozoic and Precambrian platforms, one would not expect neither fluids nor partial melts to play an important role in velocity or attenuation anomalies in the upper 100 km. Except for special tectonic settings (e.g. Nolet & Zielhuis 1994),  $Q_s$  and  $V_s$  anomalies there are likely to be compositional. In tectonically active regions, even at near-solidus temperatures, the dependence of velocity and attenuation on temperature variations becomes strongly non-linear (e.g. Sato & Sacks 1989; Sato *et al.* 1989; Jackson *et al.* 1992). A small percentage partial melting dramatically affects shear modulus and shear wave attenuation through direct thermoelastic losses (Anderson 1989; Sato *et al.* 1989; Jackson *et al.* 1992; Karato 1993). Because the effect

of  $T$  on  $Q_s$  and  $V_s$  is different (a sharp increase in  $Q_s$  occurs at lower temperature than a drop in  $V_s$  associated with partial melting, Fig. 1), this can decrease global correlations between  $V_s$ ,  $Q_s$  and  $T$ . Romanowicz (1995) argues that exponential dependence of attenuation on  $T$  favours hotter anomalies (even when cold anomalies have larger amplitude), while velocity is more sensitive to cold anomalies as a result of its linear dependence on  $T$ . A presence of volatiles or fluids associated with subduction zones (Carlson & Miller 2003) can significantly lower the solidus and thus have a strong effect on  $Q_s$  and  $V_s$ . Thus, we conclude that the local presence of partial melts and/or fluids may weaken global  $V_s$ – $Q_s$ – $T$  correlations.

Because we are unable to separate the effects of the lack of resolution in our data from true non-thermal effects, we next examine regional qualitative correlations between  $V_s$ ,  $Q_s$  and  $T$  (Section 4) to test if, regardless of low global correlation values, temperature plays a governing role in producing regional seismic anomalies. We further use theoretical and laboratory data to quantify which part of observed variations in seismic attenuation and shear wave velocities can be predicted from thermal anomalies alone (Section 5).

#### 4 QUALITATIVE COMPARISON OF $V_s$ , $Q_s$ AND $T$ ANOMALIES IN THE UPPER MANTLE

##### 4.1 General remarks

We visually compare maps (Figs 2–4) of seismic and temperature anomalies in the upper mantle to examine if there is a qualitative regional correlation between the parameters. Because at 50 km depth the crustal structure has the largest influence on the  $V_s$  and  $Q_s$  values, we place greater emphasis on the maps for 100 and 150 km depth, with particular focus on the results for 100 km depth, where the resolution of both of the seismic models is the highest. For a typical continental geotherm (corresponding to surface heat flow of *ca* 60 mW m<sup>-2</sup>), lithospheric temperatures at 100 km depth are close to 1000 °C. At seismic frequencies, a strong solid-state viscoelastic relaxation expressed as a sharp increase in seismic attenuation is expected in the upper-mantle crystalline rocks at high homologous temperatures (e.g. Berckhemer *et al.* 1982). This means that anelastic anomalies of thermal origin should be best correlated with temperature at a depth of 100 km. Indeed, white areas with a near-zero  $Q_s$  anomaly generally correlate with a 1000–1100 °C isotherm (Fig. 3).

A comparison of the three maps presented in Figs 2–4 reveals the existence of a strong overall qualitative correlation between the signs of  $V_s$ ,  $Q_s$  and  $T$  anomalies for most of the continental lithosphere at depths of 50–150 km. At first look, there is less correlation between the three maps for a depth of 150 km. The eye catches strong positive seismic anomalies over North America and Eurasia, where temperature anomalies appear to be much weaker. Nevertheless, the maps reveal a good qualitative correlation between the signs of the anomalies at this depth as well.

##### 4.2 Regions where all three parameters correlate

There are two types of regions where all three anomalies correlate: cold regions, which include most of the Precambrian cratons, and hot regions such as continental rifts, Cenozoic orogens and, unexpectedly, the Sino-Korean craton. The signs of all three anomalies correlate within both of the groups throughout the entire lithosphere,

which suggests a primarily thermal origin of most of the seismic anomalies.

Hot (red) anomalies (low  $V_s$ , low  $Q_s$  and high  $T$ ) are found in all three maps (Figs 2–4) in most tectonically active regions of the continents. At all depths, blue, cold anomalies with high  $V_s$ , high  $Q_s$  and low  $T$  dominate in cratonic regions (see Fig. 5 for a location map). For example, high  $V_s$ , high  $Q_s$  and low  $T$  are seen beneath the West African craton and the western-central part of Australia, in agreement with recent regional shear wave tomography studies (Simons *et al.* 1999; Debayle & Kennett 2000).

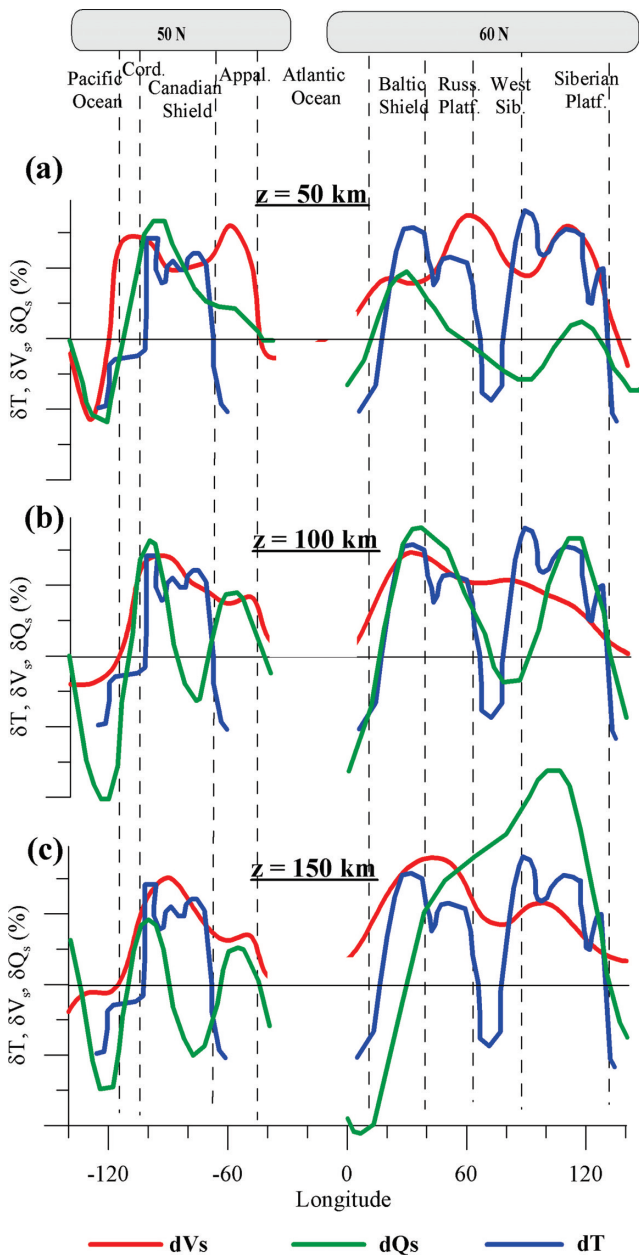
On the whole, we find a better agreement between  $T$  and  $Q_s$  for large cratons (the Canadian shield, the Baltic shield, the East European platform and the Siberian craton) than for small ones. For example, a strong correlation between all three anomalies ( $V_s$ ,  $Q_s$  and  $T$ ) is observed at 100 km depth beneath most of Eurasia, with the strongest anomalies beneath the East European and the Siberian cratons separated by a hot high-temperature and high-attenuation anomaly beneath the Palaeozoic West Siberian basin (WSB). However, at 150 km depth, pronounced low attenuation in the WSB (Figs 4b and 7c) does not correlate with high temperature (Fig. 4c).

In the Canadian shield, the magnitude of the  $T$  anomaly is significantly weaker than the magnitude of the  $Q_s$  and  $V_s$  anomalies, probably because all borehole heat flow data are located at the periphery of the Canadian shield. The values of the temperature anomalies in our model (Fig. 3c) are close to the estimates of Goes & van der Lee (2002), who calculated lithospheric temperature in North America from seismic tomography data, assuming that  $V_s$  anomalies are caused by temperature variations. The good agreement between two independent estimates of lithospheric temperatures beneath stable North America indicates that shear wave velocity anomalies in the Precambrian lithosphere are indeed chiefly of thermal origin.

The Early Proterozoic Sino-Korean craton shows near-zero anomalies in the three maps at 100 km depth and hot anomalies in the maps for 150 km. Surface wave tomography (Huang *et al.* 2003), heat flow data (Artemieva & Mooney 2001) and studies of mantle xenoliths (Xu *et al.* 2000) indicate an atypically thin (120–150 km) Sino-Korean lithosphere compared with other Precambrian cratons. Furthermore, the lower part (below 80–140 km depth) of the cratonic lithosphere could have been largely replaced by more fertile Phanerozoic material (Griffin *et al.* 1998). Thus, part of the  $V_s$  and  $Q_s$  variations there may probably have a non-thermal origin (possibly including compositional variations at *ca* 100 km depth and partial melting below *ca* 150 km depth).

##### 4.3 Regions where only seismic anomalies ( $V_s$ and $Q_s$ ) correlate

Of particular interest are the regions where the seismic anomalies  $V_s$  and  $Q_s$  have the same sign, but do not correlate with temperature anomalies, because such pattern can be indicative of a clearly non-thermal origin of the seismic anomalies. The East African rift zone is one of the regions that exhibits such non-correlated behaviour of seismic and thermal anomalies. However, this rift is too narrow to be reliably resolved by seismic models used in the present study. The Tibet–Himalayan orogen is consistently evident only in the map of  $T$ . However, temperature estimates for this tectonically active region based on regional metamorphism data should be interpreted with caution. A large time delay in the upward propagation of a thermal anomaly also implies that the correlation between seismic and thermal models should be best in stable regions that did not undergo tectonic or magmatic activity over a long geological time. Thus, we cannot identify any region where seismic anomalies are clearly of a non-thermal origin.



**Figure 7.** Comparison of  $Q_s$ ,  $V_s$  and  $T$  at a depths of 50, 100 and 150 km along the latitude 50°N in North America and the latitude 60°N in Eurasia. The values of  $V_s$ ,  $Q_s$  and  $T$  are averaged over a 10° band, centered correspondingly on 50°N or on 60°N. Large ticks on the y-axis correspond to the 4 per cent anomaly for  $V_s$ , 40 per cent anomaly for  $Q_s$  and -20 per cent anomaly for  $T$ .  $Q_s$  and  $V_s$  anomalies are shown as variation with respect to PREM.  $T$  anomaly is shown with respect to a typical continental geotherm of 55–60 mW m<sup>-2</sup> (i.e. ca 750, 1000 and 1250°C at three depths, correspondingly). An accuracy in the amplitudes of  $T$  model is ca 10 per cent, so that the signs of the anomalies are the robust features. The lateral resolution of thermal and seismic models is different (ca 20° for  $Q_s$  and  $V_s$  anomalies and ca 10° for  $T$  anomalies), such that their peaks may not necessarily coincide. The resolution of seismic data is the lowest at 50 km depth as a result of crustal effects. On the whole, there is a good general agreement between  $V_s$  and  $T$  anomalies. However, at 100 km depth where the resolution of the attenuation model is the highest,  $Q_s$  and  $T$  anomalies in Eurasia show an excellent agreement.

#### 4.4 Regions where $V_s$ and $T$ anomalies do not correlate with $Q_s$

The regions with non-correlated seismic anomalies in the lithosphere (which include the SW Baltic shield and the South African craton) deserve special attention, as both  $V_s$  and  $Q_s$  anomalies used in this study were derived in a joint inversion of the same carefully selected data set (Billien 1999; Billien *et al.* 2000). According to experimental studies on polycrystalline olivine at seismic frequencies (e.g. Jackson 2000), attenuation in shallow lithosphere, however, needs not correlate with temperature and shear wave velocities, because at low  $T$  (< 600°C) the upper-mantle rocks exhibit an almost elastic behaviour, while the values of  $Q^{-1}$  are essentially scattered.

The late Proterozoic southwestern part of the Baltic shield is close to the transition zone from depleted Precambrian lithosphere to fertile, younger lithosphere. The transition (the Trans-European suture zone, TESZ) is marked by a sharp contrast in crustal structure, surface heat flow and lithospheric thickness. In maps for 50 and 100 km depths, the  $V_s$ ,  $Q_s$  and  $T$  anomalies approximately follow the NW–SE orientation of TESZ. However, at 150 km depth the attenuation anomaly has a N–S orientation (Fig. 4b). This anomaly may imply the presence of partial melt and/or fluid inclusions in the metasomatized lithosphere, which at near-solidus temperatures would affect  $Q_s$  stronger than  $V_s$  (Mitchell 1995). This hypothesis is supported by electrical studies (Jones 1984, 1999), which imply a presence of a well-conducting asthenosphere at ca 160–190 km depth in the western part of the Baltic shield. Deep mantle reflections seen in the BABEL seismic profile in the same region give evidence for the Proterozoic subduction (BABEL Working Group 1990) and thus regional fluid enrichment of the lithosphere can be expected there. Part of the apparent anelastic attenuation anomalies in the Baltic shield may also be the result of scattering effects in the upper mantle that was reworked during the late Proterozoic tectonic activity in the region (e.g. Gaal & Gorbatshev 1987), when large volumes of rapakivi granitic intrusions were emplaced (Haapala & Rämö 1991), followed by Palaeozoic rifting (e.g. Neumann *et al.* 1995). As a result, the marginal part of the cratonic lithosphere could have been partly detached and/or metasomatized.

Surprisingly, southern Africa, including the Kaapvaal craton, does not show the expected high  $V_s$  anomaly typical for other cratons, while temperature estimates suggest that the lithosphere is cold (e.g. Jaupart & Mareschal 1999; Artemieva & Mooney 2001). Moreover, it has a strong mantle attenuation, and  $V_s$  and  $Q_s$  anomalies disagree there at all lithospheric depths. However, the discrepancy between the signs of the anomalies is likely to be a resolution artefact. The apparently highly attenuating region at 100 km depth beneath South Africa may partly result from smearing of a strong offshore anomaly (see the maps in Billien *et al.* 2000). Similar explanation of low  $V_s$  uncorrelated with low  $T$  applies to the Indian shield, where a strong offshore seismic anomaly is observed.

#### 4.5 Regions where $Q_s$ and $T$ do not correlate with $V_s$

Regions where neither  $Q_s$  nor  $T$  anomalies correlate with  $V_s$  anomalies in the deep lithosphere ( $z \sim 100$ – $150$  km) include the West Siberian Basin (WSB) and the Arabian shield. However, the lateral dimensions of the Arabian shield are too small to be resolved by the seismic models used in the present study and only a few heat flow measurements are available there. We favour a compositional origin of the velocity anomalies beneath the WSB, because laboratory and theoretical studies indicate that seismic velocities are more sensitive

to compositional variations than seismic attenuation (e.g. Bina & Silver 1997). We further address this question in Section 5.3.

#### 4.6 Comparison of $V_s$ , $Q_s$ and $T$ along a 50°N–60°N profile

A global profile at 50°N latitude in North America and at 60°N latitude in Eurasia permits the examination of regional correlations between  $V_s$ ,  $Q_s$  and  $T$  in more detail, because the temperature model used in this study is well constrained only for stable continents. Thus, our choice of a profile has been dictated by a wish to compare  $V_s$ ,  $Q_s$  and  $T$  primarily in platform regions that have good heat flow data coverage.

As discussed earlier, there is a good general agreement between  $V_s$  and  $T$  anomalies for the cratonic regions (the Canadian shield, the Russian (East European) platform and the Siberian platform). At 100 km depth (Fig. 7b) where the resolution of the attenuation model is the highest,  $Q_s$  and  $T$  anomalies in Eurasia show a perfect agreement. Increased attenuation at the southeastern margin of the Canadian shield observed at 100–150 km depth results from smearing of a high attenuation anomaly centered over the Appalachians (compare with Fig. 3b). As discussed in the previous section, the non-correlated behaviour of the anomalies in the WSB is likely to be caused by compositional variations. Thus, we conclude that, despite relatively low values of correlation coefficients, a significant part of seismic anomalies in stable continental lithosphere has a thermal origin.

### 5 COMPARISON OF SEISMICALLY MEASURED $V_s$ AND $Q_s$ AND THEORETICAL $V_s^T$ AND $Q_p^T$

#### 5.1 Calculation of theoretical $Q_p^T$ in the upper mantle from $T$

With the available data set we were unable to confirm that, statistically,  $T$  is the dominant factor for  $Q_s$  and  $V_s$  variations on a global scale (Section 3). This appears to contradict numerous laboratory measurements that demonstrate the temperature dependence of seismic attenuation and provide the basis for estimates of attenuation in the upper mantle of a purely thermal origin,  $Q_p^T$  (eq. 1). Moreover, qualitative comparison shows that on the regional scale the signs of seismic and temperature anomalies in the continental lithosphere are correlated (section 4). Though the robust extrapolation of experimental data to seismic attenuation in the real Earth is controversial (e.g. Jackson 2000), a comparison of  $Q_p$  calculated from  $Q_s$  from global anelastic tomography with theoretical predictions for  $Q_p^T$  from temperatures would permit to distinguish attenuation anomalies in the upper mantle that have a clearly thermal origin.

Sato & Sacks (1989) and Sato *et al.* (1989) have measured the attenuation in dry peridotite at ultrasonic frequencies (60–900 KHz) in the temperature range  $0.95 < T_m/T < 1.17$  (where  $T_m$  is solidus temperature) and at pressures of 200–730 MPa. Their assumption of a dry upper-mantle composition is in agreement with recent studies of the electrical conductivity of the stable continental mantle (Hirth *et al.* 2000). The linear dependence of  $Q_p$  on the homologous temperature  $T_m/T$  (at ultrasonic frequencies) allows extrapolation of the laboratory results to higher pressures through the pressure-dependent solidus temperature  $T_m(P)$  (in Kelvin):

$$Q_p = Q_{pm} \exp\{g[T_m(P)/T - a]\}, \quad (5)$$

where  $Q_{pm} = 3.5 + P/0.073$  (here  $P$  is pressure in GPa),  $a$  and  $g$  are experimentally determined parameters ( $a$  is close to 1, while

$g$  is strongly  $T$  dependent and increases from 6.75 to 13.3 with an increase of homologous temperature) (Sato *et al.* 1989).

The measurements of Sato *et al.* (1989) were done at ultrasonic frequencies, while mantle attenuation is believed to be strongly frequency dependent (Karato 1998). Romanowicz & Durek (2000) argue that the procedure suggested by Sato *et al.* (1989) to extrapolate their relationship (eq. 5) between  $Q_p$  and  $T_m/T$  to the seismic frequencies may not be valid, as mechanisms of seismic attenuation other than the grain boundary process assumed by Sato & Sacks (1989) (e.g. dislocation and microcreep processes) can be equally important at seismic conditions (Karato 1998). Nevertheless, eq. (5) permits the making of rough estimates of  $Q_p$  from known geotherms, but only in the vicinity of solidus temperatures (*ca* 1200 to 1500 °C). Thus, the approach is not valid for calculations of  $Q_p^T$  in the shallow lithosphere or in the cold lithosphere of Precambrian cratons.

Experimental data of Jackson *et al.* (1992) provide basis for estimation of seismic attenuation in a wide temperature range. Their laboratory measurements of attenuation at seismic frequencies (0.01–1 Hz) were made on dry dunite at pressures up to 300 MPa and at temperatures of 20, 600, 800 and 1000 °C. They show that  $Q_p$  strongly depends on the oscillation period,  $\tau$ , and between 600 and 1000 °C increases approximately exponentially with increasing temperature  $T$ :

$$1/Q_p = 1/Q_{p0} \tau^\alpha, \quad (6)$$

where  $1/Q_{p0}$  (i.e. the value of  $1/Q_p$  at a 1 s period) exponentially depends on temperature, while the coefficient  $\alpha$  is almost temperature independent and may range from 0.15 to 0.30 for different values of  $\tau$ ,  $T$  and  $P$ . We have approximated the experimental data from Jackson *et al.* (1992) as ( $T$  in °C)

$$1/Q_{p0} = \exp(0.00267 * T - 6.765) \quad (7)$$

and have taken  $\alpha \sim 0.20$ – $0.25$ .

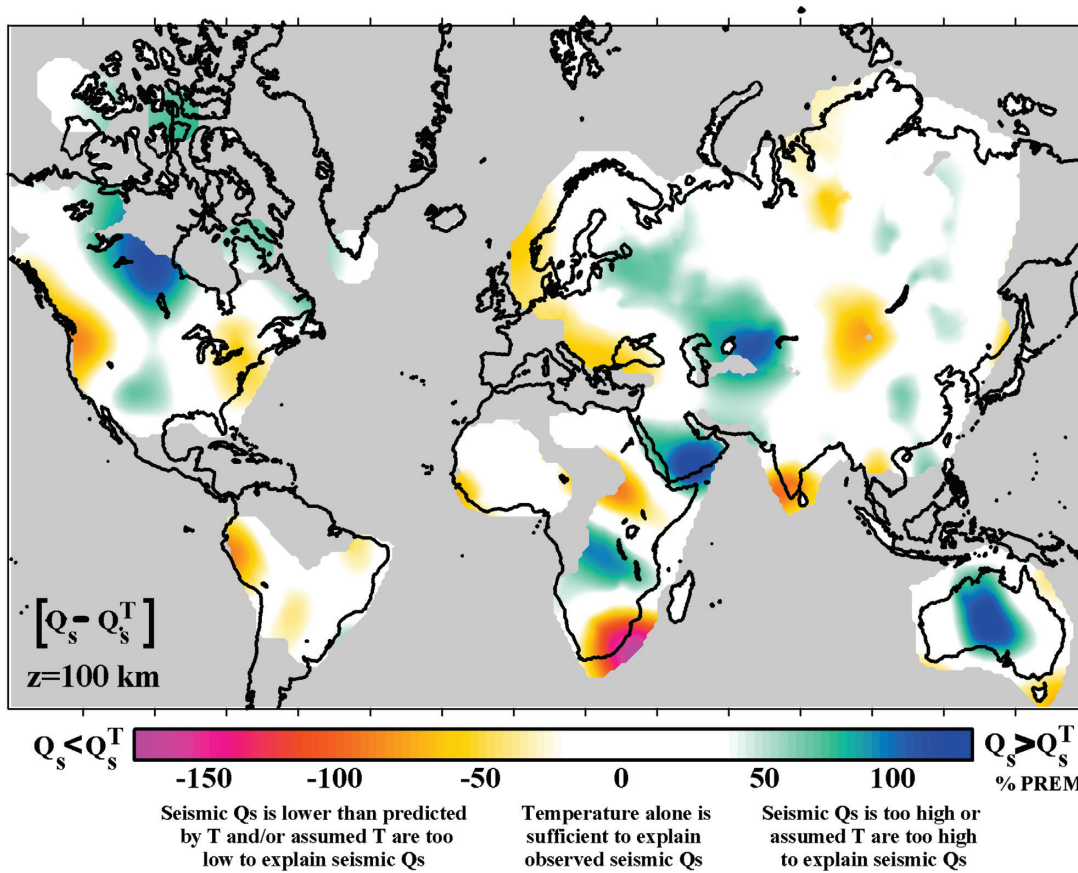
Because the attenuation model based on seismic inversion has the highest resolution at 100 km depth, we focus our analysis at this depth in the upper mantle, for which we compare seismic  $Q_s$  with theoretical estimates. We use eqs (6–7) to calculate theoretical values of  $Q_p^T$  in the continental upper mantle from temperatures (Artemieva & Mooney 2001) (Fig. 3c).  $P$ -wave attenuation,  $Q_p^T$ , at a depth of 100 km beneath the continents is calculated from the data on lithospheric temperatures (Artemieva & Mooney 2001) for an oscillation period  $\tau = 40$  s and then is converted into  $Q_s^T$  values, following the relationship (e.g. Anderson *et al.* 1965)

$$Q_p^{-1} = 4/3 V_s^2 / V_p^2 Q_s^{-1} \quad (8)$$

and assuming that for the upper-mantle velocities  $Q_p = 2.25 Q_s$ . Because the seismic  $Q_s$  model was calculated with respect to the ocean-dominated PREM model, there is a systematic shift of the absolute  $Q_s$  values on continents. Thus, the zero anomaly in Fig. 8 (which shows the difference between the seismic  $Q_s$  (Fig. 3b) and the  $Q_s^T$  values) is centered over the average difference between seismic  $Q_s$  and theoretical  $Q_s^T$  values, i.e. over  $\delta Q_s = Q_s - Q_s^T \sim 80$ .

#### 5.2 Calculation of theoretical $V_s^T$ at 100 km depth from $T$

Laboratory measurements of the temperature dependence of density, elastic parameters and, sometimes, direct measurements of the seismic velocities in mantle rocks and rock-forming minerals (Berckhemer *et al.* 1982; Kampfmann & Berckhemer 1985; Sato *et al.* 1989; Jackson *et al.* 1990) provide the basis for estimates of seismic velocities in the upper mantle of a purely thermal origin. A forward



**Figure 8.** Map of the difference  $\delta Q_s$  between (i)  $Q_s$  at 100 km depth as calculated from surface waves (Billien *et al.* 2000) and (ii) the value of  $Q_s^T$  estimated from experimental data of Jackson *et al.* (1992) for  $\tau = 40$  s and the lithospheric temperatures (Artemieva & Mooney 2001). A large part of the continental lithosphere with diverse ages and tectonic settings fits a purely thermal model for seismic  $Q_s$  (white areas). Many active regions appear as more attenuating ( $\delta Q_s < 0$ ), probably as the result of the presence of fluids and partial melts. The anomalies over South Africa and India can be produced by smearing of strong offshore attenuation anomalies.  $Q_s$  anomalies in some cratonic regions (central Australia and the Canadian shield) are higher than predicted by  $Q_s^T$  models ( $\delta Q_s > 0$ ), suggesting that lithospheric temperatures can be somewhat lower than suggested in Fig. 3(c). Alternative explanation includes scattering on large-scale structural heterogeneities in the cratonic lithosphere, which was formed by collisions of numerous continental and oceanic terranes.

calculation of theoretical  $V_s^T$  values from temperature data is free of many uncertainties of inverse modelling of temperatures from seismic  $V_s$  because it is not necessary to quantify the effects of composition or the presence of partial melts and fluids on seismic velocities (e.g. Berckhemer *et al.* 1982; Schmeling 1985; Popp & Kern 1993). To calculate theoretical velocities from  $T$  (Fig. 3c), we first assume a linear dependence of velocity on temperature with  $\partial V_s / \partial T = 0.35 \text{ m s}^{-1} \text{ K}^{-1}$  (Sumino & Anderson 1982) and for a 100 km depth calculate  $V_s^T$  (elastic) as velocity anomaly  $\partial V_s$  with respect to the PREM value, produced by a temperature variation  $\partial T$  with respect to a typical continental geotherm of  $55\text{--}60 \text{ mW m}^{-2}$  (i.e. *ca*  $1000^\circ\text{C}$  at  $z = 100 \text{ km}$ ).

In the Earth's mantle, anelasticity significantly affects seismic velocities at seismic frequencies (Karato 1993) and, if not accounted for, can result in overestimated  $V_s^T$  anomalies. Because the correlation coefficient between  $Q_s$  and  $V_s$  calculated in seismic inversion is 0.66 at  $z = 100 \text{ km}$  (Section 3), probably partly as a result of the effect of scattering on  $Q_s$ , we account for anelasticity as (Minster & Anderson 1981):

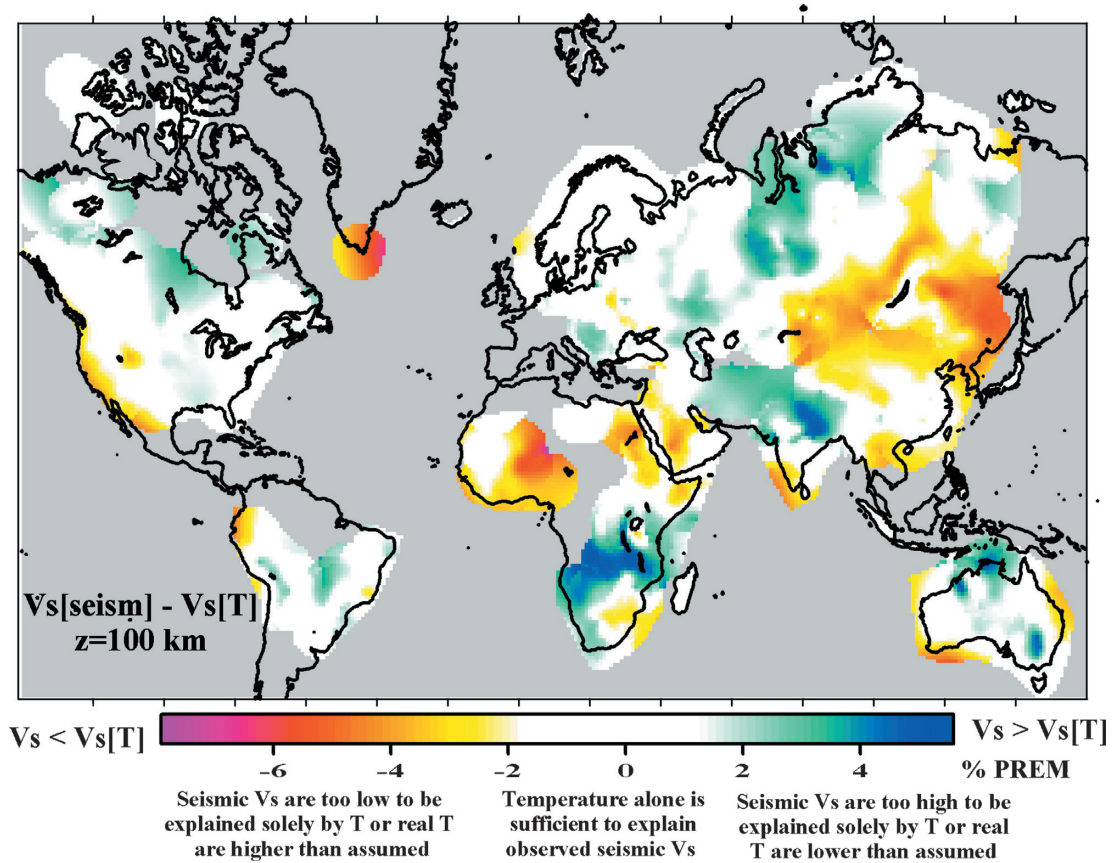
$$V_s^T = V_s^T(\text{elastic})[1 - CQ^{-1}(\tau, P, T)], \tag{9}$$

where  $C = 2/\tan(\pi\alpha/2)$ ,  $Q^{-1}(\tau, P, T)$  is given by eq. (1) and the numerical values of the parameters are  $\alpha = 0.2$ ,  $E^* = 560 \text{ kJ mol}^{-1}$  at  $z = 100 \text{ km}$ , and the oscillation period  $\tau = 40 \text{ s}$ . Following

Sobolev *et al.* (1996), who calibrated the scaling constant  $A$  (eq. 1) by fitting average attenuation in the upper mantle by the old ocean geotherm, we take  $A = 0.148$ .

### 5.3 Which part of observed variations in seismic $Q_s$ and $V_s$ can be predicted from thermal anomalies alone?

Figs 8–9 show differences  $\delta Q_s$  and  $\delta V_s$  between seismic  $Q_s$  and  $V_s$  (Figs 3a and b) and theoretical  $Q_s^T$  and  $V_s^T$  calculated from temperatures (Fig. 3c). Small-scale anomalies result from a different lateral resolution of seismic and temperature models and should be interpreted with caution. Figs 8–9 indicate that in *ca* 50 per cent of the continents seismic attenuation and velocity anomalies can be explained by thermal effects alone (white areas). However, these regions do not show any systematic correlation either with the tectonic setting, or with geological age. Tables 2(a) and (b) provide a summary of the continental regions, where  $Q_s$  and  $V_s$  anomalies calculated from seismic inversion significantly differ from  $Q_s^T$  and  $V_s^T$  expected from temperature variations alone. We estimated how much higher (or lower) mantle temperatures should have been to explain the observed  $\delta Q_s = Q_s - Q_s^T$  and  $\delta V_s = V_s - V_s^T$  misfits. Fig. 10 illustrates how  $\delta Q_s$  and  $\delta V_s$  anomalies were converted into  $\delta T$  anomalies: however, there is a large uncertainty in a conversion of  $\delta Q_s$  anomalies into  $\delta T$ .



**Figure 9.** Map of the difference  $\delta V_s$  between (i)  $V_s$  at 100 km depth as calculated from surface waves (Billien 1999) and (ii) the value of  $V_s^T$  estimated from the lithospheric temperatures (Artemieva & Mooney 2001) with account for anelasticity, for  $\tau = 40$  s. Only small areas of the continental lithosphere fit a purely thermal model for seismic  $V_s$  (white areas). In many active regions  $\delta V_s$  is ca -3 per cent (which corresponds to ca 50–70 per cent of the total  $V_s$  anomaly), probably as a result of the presence of fluids and partial melts. Positive  $\delta V_s$  anomalies in most of the cratonic regions can be explained by Fe depletion of lithospheric keels. Several Phanerozoic regions have  $\delta V_s$  too large (ca 3 per cent) to be explained by overestimated  $T$ , these anomalies seem to be associated with past and present subduction zones.

5.3.1 Regions with  $Q_s < Q_s^T$  and  $V_s < V_s^T$

Red colors in both of the maps (Figs 8–9) outline regions where seismic anomalies could have been explained if upper-mantle  $T$  were higher than assumed (Fig. 3c). Surprisingly, there is not much of overlap between the regions with  $Q_s < Q_s^T$  and  $V_s < V_s^T$  (Table 2a). Areas where seismic  $V_s$  are significantly lower than predicted from  $T$  (Fig. 9) include mostly tectonically active regions, but negative  $\delta Q_s$  anomalies fall both within active and stable continental regions, suggesting that no single mechanism gives rise to these anomalies. The Arabian shield is the only region where the signs of  $\delta Q_s$  and  $\delta V_s$  anomalies disagree, probably because of a lack of lateral resolution.

Three cratonic regions show negative  $\delta Q_s$  anomalies (Table 2a): however, theoretical velocities well agree there with seismic  $V_s$ . Thus, we suggest that seismic  $Q_s$  anomalies over the South African and Indian cratons are smearing artefacts of strong offshore attenuation anomalies. We are less certain about the origin of the  $\delta Q_s$  anomaly over a part of the Canadian shield: however, its low amplitude suggests that temperature plays an important role in producing seismic anomalies in all three of these cratonic regions.

The only regions that are red in both of the maps ( $\delta Q_s < 0$  and  $\delta V_s < 0$ ) are the western USA and the mountain country of central Asia (Tien-Shan, Sayans and Altai mountains, and Inner Mongolia). These tectonically active regions have a transient thermal regime

within the upper mantle and thus temperatures cannot be constrained from a steady-state thermal conductivity equation. Petrologic and non-steady-state geothermal constraints available for these regions (e.g. Lachenbruch & Sass 1977; Polyakov *et al.* 1988; Henry *et al.* 1997; Le Pichon *et al.* 1997) provide constraints on mantle temperatures, but an uncertainty in  $T$  values in these regions is much larger than for stable continents. Estimates show that  $\delta Q_s$  and  $\delta V_s$  could be explained if mantle  $T$  were above the mantle adiabat, implying that partial melts significantly decrease both  $Q_s$  and  $V_s$ . This conclusion is supported by regional seismic studies in western USA and Tien Shan (Humphreys & Dueker 1994; Roecker *et al.* 1993) that indicate low seismic velocities in the uppermost mantle and suggest intensive melting at ca 100 km depth. Thus, we conclude that all continental regions with reliable constraints on mantle temperatures, where  $Q_s > Q_s^T$  and  $V_s > V_s^T$ , require the presence of partial melt (and/or fluids) in the uppermost mantle to explain observed seismic anomalies.

5.3.2 Regions with  $Q_s > Q_s^T$  and  $V_s > V_s^T$

There is a striking similarity between the regions with  $Q_s > Q_s^T$  and  $V_s > V_s^T$  anomalies (Table 2b). Regions where both  $Q_s$  and  $V_s$  calculated from seismic inversion are higher than expected from

**Table 2a.** Regions (Figs 8–9) where  $V_s$  and  $Q_s$  anomalies suggest mantle temperatures higher than assumed (alternative and preferred explanation is a presence of partial melts/fluids).

Regions and $T$ at $z = 100$ km	$\delta V_s = V_s - V_s^T$ (per cent PREM) For $V_s^T$ with account for anelasticity <sup>(a)</sup>	$\delta T$ (°C) required to explain $\delta V_s$		Regions and $T$ at $z = 100$ km	$\delta Q_s = Q_s - Q_s^T$ (per cent PREM)	$\delta T$ (°C) required to explain $\delta Q_s$
		For $\delta V_s$ with account for anelasticity <sup>(a)</sup>	For $\delta V_s$ from linear $V_s$ - $T$ relationship <sup>(b)</sup>			
Tien Shan, Altai, Sayans and Inner Mongolia (°) (1100 °C)	-3–4 per cent	+250 °C	+300 °C	Tien Shan, Altai, Sayans and Tarim (°) (1100 °C)	-50–100	Non-defined, in excess of +600 °C
Western North America (*) (°) (1300 °C)	-2–3 per cent	+150 +180 °C	+300 °C	Western North America (**) (°) (1300 °C)	-50–100	Already above adiabat
S Greenland (900 °C)	-4–5 per cent	+250 +300 °C	+500 °C	Abitibi province (Canadian shield) and central part of Grenville province (800–1000 °C)	-50	Non-defined, in excess of +600 °C
Cameroon volcanic line (W Africa) (800 °C) (°)	-3–4 per cent	+180 +250 °C	+500 °C	South Africa (***) (800–1100 °C)	-150	Non-defined, in excess of +600 °C
Arabian shield and Nubian shield (900–1100 °C)	-3 per cent	+180 °C	+300 +400 °C	Indian shield (****) (800–900 °C)	-150	Non-defined, in excess of +600 °C
S part of the Russian Far East (1100 °C) (°)	-4–5 per cent	+250 +300 °C	+500 °C	Western part of the Siberian craton (500–700 °C)	-50	+200 +400 °C
				Trans-European Suture Zone (900–1100 °C)	-50	Non-defined, in excess of +600 °C

<sup>(a)</sup> For  $V_s^T$ ,  $\delta V_s$  and  $\delta T$  calculated after Deschamps *et al.* (2002).

<sup>(b)</sup> For  $V_s^T$ ,  $\delta V_s$  and  $\delta T$  calculated from  $\partial V_s / \partial T = 0.35 \text{ m s}^{-1} \text{ K}^{-1}$ .

<sup>\*</sup>  $T$  taken from published petrological and non-steady-state constraints on the thermal regime. No heat flow data,  $T$  at the maps is interpolation artefact.

<sup>\*</sup>  $V_s$ -to- $T$  conversion used to constrain Fig. 9 can be invalid for this region because of high homologous  $T$ .

<sup>\*\*</sup> Eq. (6) used to constrain Fig. 8 can be invalid for this region because of high homologous  $T$ .

<sup>\*\*\*</sup> The anomaly most likely is an extension of a strong offshore  $Q_s$  anomaly.

**Table 2b.** Regions (Figs 8–9) where  $V_s$  and  $Q_s$  anomalies suggest mantle temperatures lower than assumed (alternative and preferred explanation is Fe depletion of cratonic lithosphere).

Regions and $T$ at $z = 100$ km	$\delta V_s = V_s - V_s^T$ (per cent PREM) For $V_s^T$ w/unt for anelasticity <sup>(a)</sup>	$\delta T$ (°C) required to explain $\delta V_s$		Regions and $T$ at $z = 100$ km	$\delta Q_s = Q_s - Q_s^T$ (per cent PREM)	$\delta T$ (°C) required to explain $\delta Q_s$
		For $\delta V_s$ with account for anelasticity <sup>(a)</sup>	For $\delta V_s$ from linear $V_s-T$ relationship <sup>(b)</sup>			
East European platform (750–850 °C)	+1 +2 per cent	-70 – 150°C	-100 – 250°C	East European platform (750–850 °C)	+50	-250°C
North and Central Australia (900–1000 °C)	+1 +2 per cent	-70 – 150°C	-100 – 250°C	Central Australia (900–1000 °C)	+150	In excess of -600°C
Slave and Hearne provinces	+1 +2 per cent	-70 – 150°C	-100 – 250°C	Slave and Hearne Provinces (Canadian shield) (900 °C)	+100	-350 – 450°C
Canadian shield (900 °C)	+1 per cent	-70°C	-100°C	South-central USA (900–1000 °C)	+50	-250 – 300°C
Mobile belts of South Africa and Central Africa (*) (1200–1300 °C)	+5 per cent	-450°C	-400 – 600°C	Central Africa (**) (1200–1300 °C)	+100	In excess of -600°C
Himalayas and the Tethys belt (*) (1300 °C)	+3 +5 per cent	-250 – 450°C	-250 – 600°C	Himalayas and the Tethys belt (*) (1300 °C)	+50 +100	In excess of -600°C
West Siberia basin and Taimyr peninsula (*) (1100–1300 °C)	+3 +4 per cent	-250 – 350°C	-500 – 600°C	Arabian shield (900–1100 °C)	+150	In excess of -600°C

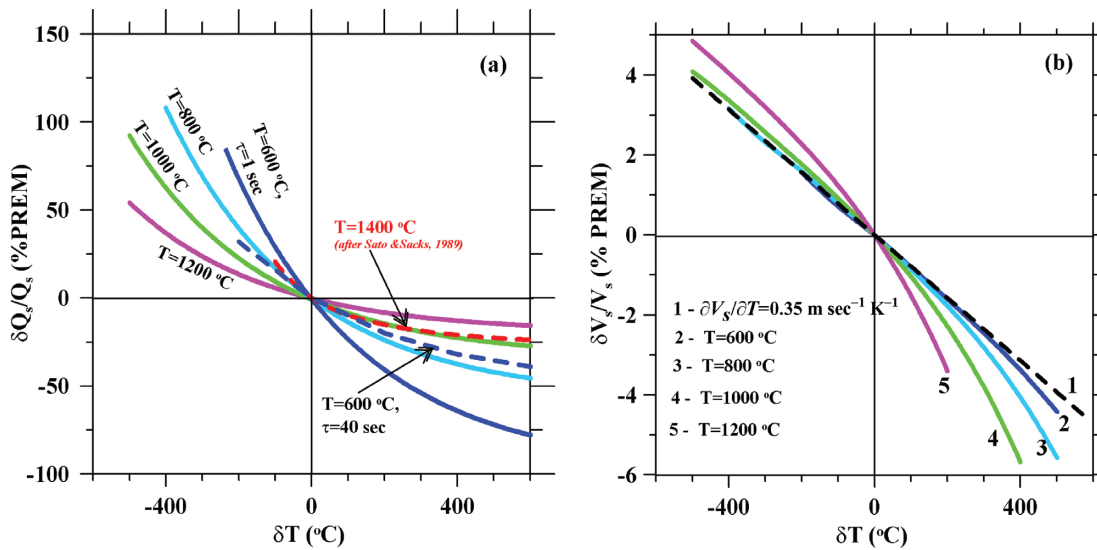
<sup>(a)</sup> For  $V_s^T$ ,  $\delta V_s$  and  $\delta T$  calculated after Deschamps *et al.* (2002).

<sup>(b)</sup> For  $V_s^T$ ,  $\delta V_s$  and  $\delta T$  calculated from  $\partial V_s / \partial T = 0.35 \text{ m s}^{-1} \text{ K}^{-1}$ .

$T$  taken from published petrological and non-steady-state constraints on the thermal regime.

\*  $V_s - T$  conversion used to constrain Fig. 9 can be invalid for this region because of high homologous  $T$ .

\*\* Eq. (6) used to constrain Fig. 8 can be invalid for this region because of high homologous  $T$ .



**Figure 10.** Relative anomalies of  $Q_s$  (a) and  $V_s$  (b) as a function of  $T$  anomalies for different upper-mantle temperatures. (a) Solid lines:  $\delta Q_s$  anomalies based on experimental data of Jackson *et al.* (1992), eqs (6–8), for  $\tau = 1$  s. For a comparison, the results for  $T = 600$  °C are also presented for  $\tau = 40$  s (dashed line). Dashed red line:  $\delta Q_s$  calculated for high homologous mantle  $T$  based on experimental data of Sato & Sacks (1989), eq. (5). (b) Solid lines:  $\delta V_s$  anomalies based on eqs (1) and (9), for  $\tau = 40$  s. For a comparison, dashed line is  $\delta V_s$  based on a linear relationship  $\partial V_s / \partial T = 0.35 \text{ m s}^{-1} \text{ K}^{-1}$  (Sumino & Anderson 1982).

the conversion of temperature anomalies (shown in blue in Figs 8–9) mostly include Precambrian terranes. Possible explanations for positive  $\delta Q_s$  and  $\delta V_s$  anomalies include overestimated lithospheric temperatures and/or unusually dry conditions within the cratonic lithosphere and/or compositional anomalies. It is surprising that the misfit  $\delta V_s$  between theoretical and seismic  $V_s$  is the same in all cratonic regions, *ca* 1–2 per cent. It can be explained by *ca* 70–150 °C lower temperatures at  $z = 100$  km than in the  $T$  model (Fig. 3c). Alternatively, a typical difference in mg# (MgO/[MgO+FeO]) between the Archean and the Phanerozoic lithosphere (0.93 and 0.89, correspondingly, e.g. Griffin *et al.* 1999) also produces an *ca* 1 per cent velocity anomaly (Deschamps *et al.* 2002). The remaining  $\delta V_s < 1$  per cent is within the accuracy of the temperature model. Because it is more uncertain to quantify  $\delta T$  required to explain a  $\delta Q_s$  misfit and the effect of depletion on attenuation is not well known, we are inclined to conclude that  $\delta Q_s$  and  $\delta V_s$  anomalies in the cratonic lithosphere are of a compositional origin (e.g. depletion in basaltic components). This conclusion is supported by recent gravity and buoyancy modelling (e.g. Artemieva 2003; Kaban *et al.* 2003). Given that cratons were formed by collisions of numerous continental and oceanic terranes, their structure can be highly heterogeneous on the lithospheric scale and thus can produce scattering even at frequencies used in the present study, resulting in  $Q_s$  and  $\delta Q_s$  anomalies.

A 1 per cent  $\delta V_s$  anomaly in southern USA (southern Great Plain) is within the accuracy of the thermal model: however, it is unlikely that upper-mantle temperatures there are lower than assumed (Fig. 3c). Besides, the region also shows a significant  $\delta Q_s$  anomaly. Recent gravity modelling for North America indicates upper-mantle compositional anomaly in approximately the same region (Mooney & Kaban 2004). Palaeotectonic reconstructions demonstrate that this region has been affected by Palaeozoic subduction (Kluth & Coney 1981). Thus, a possible explanation for  $\delta Q_s$  and  $\delta V_s$  anomalies includes compositional variations in the upper mantle (or fluids).

Other regions where seismic velocities and attenuation cannot be adequately predicted from temperatures include the Himalayas and the Tethys belt in Eurasia and south-central Africa. However, in these regions  $\delta Q_s$  and  $\delta V_s$  anomalies can be biased by the lack of reliable constraints on mantle temperatures. Alternative explanation includes the presence of strong compositional heterogeneities associated with past and present subduction zones in the regions.

The WSB is a region with one of the strongest  $\delta V_s$  anomalies (*ca* 3 per cent), which can be explained only if mantle temperature at  $z = 100$  km were *ca* 250 °C lower (i.e. if they were approximately the same as in the Precambrian cratons). Little is known about the ages of the basement rocks in the WSB. However, the entire basin is underlain by a thick layer of basaltic lavas, which have erupted simultaneously with the Siberian traps *ca* 250 Ma (Reichow *et al.* 2002); in late Palaeozoic–Mesozoic time, the basin underwent extensive rifting (Schissel & Smail 2001). Thus, it is likely that the upper mantle of the WSB still has a transient thermal regime and the present surface heat flow does not reflect high mantle temperatures associated with Phanerozoic rifting. This means that actual upper-mantle temperatures in the WSB are likely to be higher, but not lower, than assumed (Fig. 3c). The available data does not permit the making of any further speculations on the origin of a strong  $\delta V_s$  anomaly in WSB, extending to the Taimyr peninsula where Permo-Triassic basalts are also known. However, the results of the recent study of Kaban *et al.* (2003) suggest that there are large compositional differences between the Precambrian lithospheres of the East European platform and the Siberian craton, on one hand, and the Palaeozoic WSB, on the other. Thus, a compositional origin of seismic anomalies in the lithosphere of the WSB cannot be ruled out.

In summary, all positive  $\delta Q_s$  and  $\delta V_s$  anomalies can be explained by compositional variations in the upper mantle. In the cratonic regions, these variations include an Fe depletion of lithospheric keels; in Phanerozoic regions compositional anomalies can be associated

with ancient subduction zones (central Europe, southern USA and the Tethys belt) or may have another unknown non-thermal origin (West Siberian Basin).

## 6 CONCLUSIONS

Laboratory measurements of seismic velocities and attenuation in upper-mantle rocks show a strong temperature dependence. In this study, we examine the correlations between seismic shear wave velocity,  $V_s$ , inverse attenuation,  $Q_s$ , and temperature,  $T$ , in the continental upper mantle in order to investigate the influence of temperature on  $V_s$  and  $Q_s$ .

(i) The calculated correlation coefficients,  $r$ , between  $V_s$ ,  $Q_s$  and  $T$  anomalies at 50, 100 and 150 km depths do not show a global strong correlation between any pair of parameters. At 100 km depth (where both  $V_s$  and  $Q_s$  models have high resolution),  $r = 0.66$  for  $Q_s$ - $V_s$ ,  $r = -0.56$  for  $V_s$ - $T$ , and  $r = -0.33$  for  $Q_s$ - $T$ . We discuss different physical mechanisms, as well as data modelling artefacts, that can be responsible for a deterioration of global correlations from strong dependencies expected from laboratory measurements. In particular, the poor correlation between temperature and seismic parameters at subcrustal depths may indicate insufficient resolution of temperature constraints. However, the correlation between  $V_s$  and  $Q_s$ , models derived in the simultaneous inversion of the same data set, is also relatively low. This suggests that the poor correlations can be an intrinsic property of the upper mantle. Prior studies showed that, globally, attenuation and surface heat flow values correlate in the upper 250 km with a low coefficient of  $r = 0.20$ – $0.35$  (Romanowicz 1995). We found that for the continental upper mantle the correlation between  $Q_s$  and surface heat flow is even less,  $r < -0.20$  at  $z < 150$  km. Thus, with available data sets we are unable to confirm that temperature is the dominant controlling factor for  $Q_s$  and  $V_s$  variations on the global scale in the continental upper mantle.

(ii) Because the lateral resolution of our seismic and thermal models differs greatly (*ca* 3000 km vs. 200–1000 km, respectively), this difference in resolution can be partly responsible for low numerical correlations and we visually compare the maps of  $V_s$ ,  $Q_s$  and  $T$  for 50, 100 and 150 km depths. Except for the map at 50 km depth (where crustal effects have the strongest effect on  $V_s$  and  $Q_s$ ), we find a good regional qualitative inverse correlation between seismic and thermal parameters. Furthermore, we cannot identify any continental region where both  $V_s$  and  $Q_s$  do not correlate with  $T$ , which would have indicated a clearly non-thermal origin of seismic anomalies. Globally, high  $V_s$ , high  $Q_s$  and low  $T$  correlate with cratons. Thus, in agreement with previous works, we infer that positive seismic anomalies in the cratonic lithosphere are primarily of thermal (low- $T$ ) origin. Similarly, we find the expected inverse correlations (low  $V_s$ , low  $Q_s$  and high  $T$ ) for many tectonically active regions. We hypothesize that partial melts and/or fluids in the upper mantle of active regions weaken global  $V_s$ - $T$  and  $Q_s$ - $T$  correlations. The profile along 50°N–60°N, which crosses well-constrained cratonic regions, shows a remarkable correlation between  $V_s$ ,  $Q_s$  and  $T$  anomalies. Thus, we conclude that, despite low continent-scale numerical correlations between seismic and thermal anomalies, temperature variations in the upper mantle play an important, but non-unique, role in lateral variations of  $V_s$  and  $Q_s$ .

(iii) Using laboratory data on the temperature dependence of  $V_s$  and  $Q_s$  in olivine-rich rocks, we calculate theoretical  $Q_s^T$  and  $V_s^T$  at  $z = 100$  km from known temperatures. A comparison of observed and theoretical values (i.e.  $\delta Q_s = Q_s - Q_s^T$  and  $\delta V_s = V_s - V_s^T$ ) provides an estimate of a part of seismic anomalies that cannot be

predicted from temperature variations alone. We find that  $Q_s < Q_s^T$  and  $V_s < V_s^T$  anomalies are observed chiefly in tectonically active regions with high mantle temperatures, where presence of partial melts/fluids may be the cause of very low  $V_s$  and  $Q_s$  values. However, with few exceptions, regions with  $\delta Q_s < 0$  and  $\delta V_s < 0$  anomalies do not overlap, suggesting that no single mechanism gives rise to these anomalies. Negative  $\delta Q_s$  and  $\delta V_s$  anomalies could have been explained by higher mantle temperatures. However because the assumed  $T$  values are already close to the mantle adiabat, we conclude that a presence of partial melt and/or fluids in the uppermost mantle is required to explain seismic  $V_s$  and  $Q_s$ . In particular, in active regions partial melts/volatiles may be responsible for  $\delta V_s$  anomalies of *ca* 3 per cent (i.e. > 50–70 per cent of total observed  $V_s$  anomalies). Regions with positive  $\delta Q_s$  and  $\delta V_s$  include mostly cratons. A  $\delta V_s$  anomaly of *ca* 1–2 per cent can be explained if temperatures at  $z = 100$  km were 70–150 °C lower than assumed. However, we favour compositional origin of seismic anomalies in cratonic lithosphere, because  $\delta V_s$  anomalies can be well explained by a 4 per cent Fe depletion. Several Phanerozoic regions have positive  $\delta Q_s$  and  $\delta V_s$  anomalies, with  $\delta V_s \sim 3$  per cent being too large (*ca* 50–70 per cent of  $V_s$  anomalies) to be explained by overestimated temperatures. We propose that these  $\delta Q_s$  and  $\delta V_s$  anomalies have a compositional origin and can be associated with ancient subduction zones (central Europe, southern USA, Himalayas and the Tethys belt) or have other unknown non-thermal origin (West Siberian Basin).

This study shows that even if temperature variations in the lithosphere are the main cause of seismic velocity and attenuation variations, other factors, such as presence of melts/fluids and compositional and structural variations, significantly contribute to the observed anomalies in seismic properties. Improved resolution of lithospheric  $V_s$  and  $Q_s$  is needed to further quantify the factors that control seismic anomalies in the upper mantle.

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